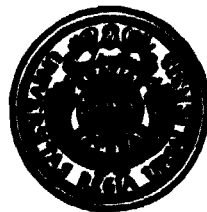


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NUCLEAR SPIN OF  $^{185}\text{Au}$   
AND HYPERFINE STRUCTURE OF  $^{186}\text{Au}$

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NUCLEAR SPIN OF  $^{185}\text{Au}$  AND HYPERFINE STRUCTURE OF  $^{188}\text{Au}$

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ABSTRACT

The nuclear spin of  $^{185}\text{Au}$ ,  $I = 5/2$ , and the hyperfine separation of  $^{188}\text{Au}$ ,  $\nu = \pm 2992(30)$  MHz, have been measured with the atomic-beam magnetic resonance method. The spin of  $^{185}\text{Au}$  indicates a deformed nuclear shape in the ground state. The small magnetic moment of  $^{188}\text{Au}$  is close in value to those of the heavier  $I = 1$  gold isotopes  $^{190}, ^{192}, ^{194}\text{Au}$ , being located in a typical transition region.

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## 1. INTRODUCTION

In a previous paper [1], we reported on spin-measurements in the gold isotopes in the mass range  $A = 186 - 189$ . These measurements were performed at ISOLDE using conventional off-line techniques. By introducing an intermittent experimental procedure, it has been possible to extend the measurements to include the nuclear spin of  $^{185}\text{Au}$  and the hyperfine structure of  $^{188}\text{Au}$ .

## 2. EXPERIMENT

The radioactive gold isotopes were obtained as daughter products of mercury produced in spallation reactions when irradiating a molten lead target of the isotope separator with 600 MeV protons from the CERN synchro-cyclotron. The ion-beam from the isotope separator was focused on the on-line oven of the atomic-beam apparatus [2]. The oven was in these measurements operated in an intermittent mode, i.e. collecting on a cold foil of molybdenum followed by a delay period to allow the gold activity to saturate, heating the oven foil to about  $1000^\circ\text{C}$  to evaporate the remaining mercury activity and finally performing the experiment, evaporating gold out of the foil at a temperature of about  $1600^\circ\text{C}$ . The different time intervals used for the experiments in  $^{185}\text{Au}$  and  $^{188}\text{Au}$  are indicated in Table 1. It is evident that only four and three exposures, respectively, could be made per hour. The activity transmitted through the atomic-beam apparatus was measured in well-shielded NaI(Tl) detectors, sensitive to the KX-rays following EC decay.

The experimental procedure for spin and moment measurements in elements with a  $J = 1/2$  electronic ground state is described in our papers on gold [1] and cesium [3]. Resonance signals corresponding to a nuclear spin  $I = 5/2$  in  $^{185}\text{Au}$  were observed at three different settings of the external magnetic field;  $\mu_B B/h = 2.060$ ,  $2.941$  and  $4.455$  MHz. No signals were observed for other half-

integer spin values up to  $1^{1/2}$ . In Fig. 1, we show the decay curve of a  $I = 5/2$  resonance signal. The half-life corresponds to that reported for the isotope  $^{185}\text{Au}$  [4],  $4.3 \pm 0.2$  min.

The nuclear spin of the 8.8 min isotope  $^{186}\text{Au}$  was measured in ref. [1] to be  $I = 1$ . In order to determine the hyperfine separation between the hyperfine levels  $F = 3/2$  and  $F = 1/2$  (c.f. Fig. 6 of ref. [3]), resonance experiments have to be made in successively stronger external magnetic fields. The single-quantum transition ( $F = 3/2, m_F = -1/2$ )  $\leftrightarrow$  ( $F = 3/2, m_F = -3/2$ ), identified by the low r.f. power required to induce it, was observed in the present experiment at two settings of the external magnetic field;  $\mu_B B/h = 35.717(11)$  MHz and  $63.566(9)$  MHz. The resonance curves, shown in Fig. 2, exhibit the same shape and width as the potassium calibration line at the two fields. The resonance frequencies  $24.236(7)$  MHz and  $43.668(9)$  MHz (errors taken to be a quarter of the line widths) were fitted to the Breit-Rabi formula (eq. 4 in ref. [3]), assuming either a positive or a negative sign of the nuclear g-factor. The resulting hyperfine separations were  $2988(26)$  MHz and  $-2996(26)$  MHz, respectively. Since there was no significant difference in the goodness of fit between the two alternatives, we give as a final result

$$\Delta\nu(^{186}\text{Au}) = \pm 2992(30) \text{ MHz}$$

The radioactive decay of the resonance samples gave a half-life of about 9 minutes, which ascertains the connection to  $^{186}\text{Au}$ .

### 3. DISCUSSION

An extensive discussion on the nuclear structure in the gold region was made in our contribution to the Cargèse Conference on Nuclei far from Stability [5]. The nuclear spin  $I = 5/2$  of  $^{185}\text{Au}$  was shown to give strong arguments for a deformed nuclear shape. The ground state may be classified by the Nilsson model configuration  $5/2 [541 3/2]$ , a configuration which appears systematically in the lighter odd-proton elements iridium, rhenium

and tantalum ( c.f. ref. [6] ). It is interesting to note that the isotope-shift measurements performed by the optical-pumping group [7] at ISOLDE have revealed a strong nuclear deformation in the isobar  $^{185}\text{Hg}$ , characterized by the neutron orbital  $[521\ 1/2]$ .

The gold isotopes  $^{188,190,192,194}\text{Au}$  all have a nuclear ground state spin  $I = 1$ . The hyperfine structure separations measured [8 - 10] in these isotopes are given in Fig. 3. In all cases only absolute values have been determined. It may be noted that the separations in  $^{188}\text{Au}$  and  $^{190}\text{Au}$  agree within the error limits, indicating similar nuclear structure of these isotopes.

In order to calculate the magnetic moments of the doubly-odd gold isotopes from the measured hyperfine constants, a direct comparison with known values [11] in the stable isotope  $^{197}\text{Au}$  may be used. Since the influence of the hyperfine anomaly might be considerable and is difficult to estimate, the magnetic moments shown in Table 2 are given without error bars. The magnetic moments are in these cases too small to allow a direct measurement, from which the hyperfine anomaly could be determined. Likewise, the uncertainty as to configuration assignments is too large to allow a theoretical calculation of the hyperfine anomaly. The nuclear ground states of these typical transition nuclei probably arise from a combination of several configurations with contributions from the low-lying  $s_{1/2}$  and  $d_{3/2}$  proton states and the  $p_{1/2}$  and  $p_{3/2}$  neutron states in the neighbouring nuclei.

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Table 1. Time sequence followed in the experiments  
on  $^{185}\text{Au}$  and  $^{188}\text{Au}$ .

Isotope	Collection on cold foil	Delay	Evaporation of mercury	Evaporation of gold
$^{185}\text{Au}$	10 min	1.5 min	1 min	3 min
$^{188}\text{Au}$	12 min	2 min	2 min	4 min



Table 2. Hyperfine separations and nuclear magnetic moments of some  $I = 1$  gold isotopes.

Isotope	Hfs separation $ \Delta\nu $ MHz	Magnetic moment $ \mu_I $ n.m. a)	Ref.
$^{188}\text{Au}$	2992(30)	0.064	b)
$^{190}\text{Au}$	3004(7)	0.064	8
$^{192}\text{Au}$	374.30(10)	0.008	9
$^{194}\text{Au}$	3489.865(32)	0.074	10

a) Because of the uncertainty as to the hyperfine anomaly, the moments are given without error bars.

b) Present work.

FIGURE CAPTIONS

Fig.1 : Decay of a resonance signal corresponding to the nuclear spin  $I = 5/2$  in  $^{185}\text{Au}$ . The signal was recorded in an external magnetic field of  $\mu_B B/h = 2.941$  MHz.

Fig.2 : Results from frequency scans in external magnetic fields of  $\mu_B B/h = 35.717(11)$  MHz and  $63.566(9)$  MHz. The resonance curves observed correspond to the single-quantum transition  $(3/2, -1/2) \leftrightarrow (3/2, -3/2)$  in the isotope  $^{188}\text{Au}$ .

Fig.3 : Hyperfine structure separations, measured in the doubly-odd gold isotopes  $^{188}, ^{190}, ^{192}, ^{194}\text{Au}$ , all having a nuclear ground state spin  $I = 1$ . The present measurement in  $^{188}\text{Au}$  is indicated by crosses, while previous results [8 - 10] are given by circles. Since the signs of the separations could not be concluded from the experiments, both possibilities are shown.

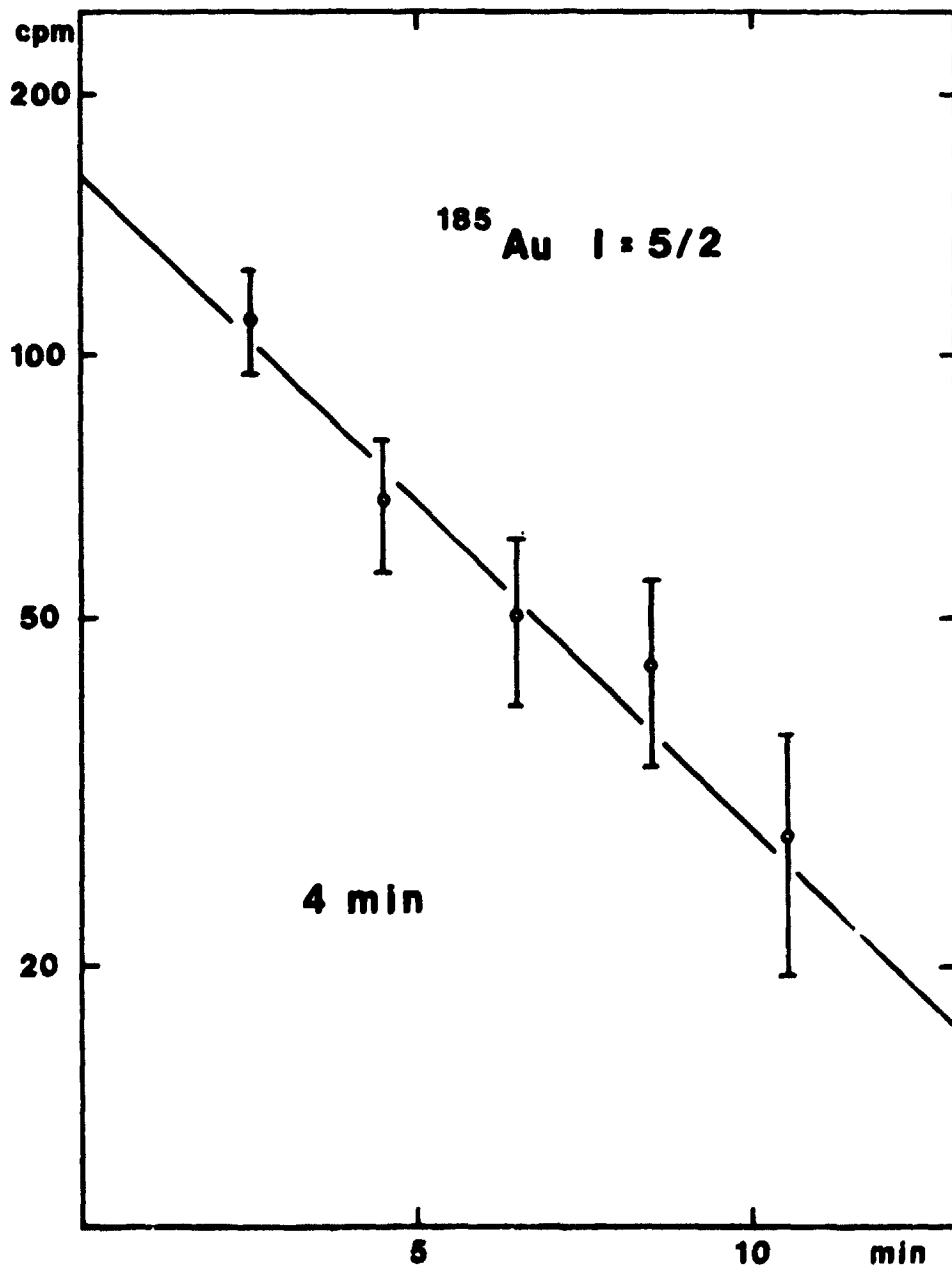


Fig. 1

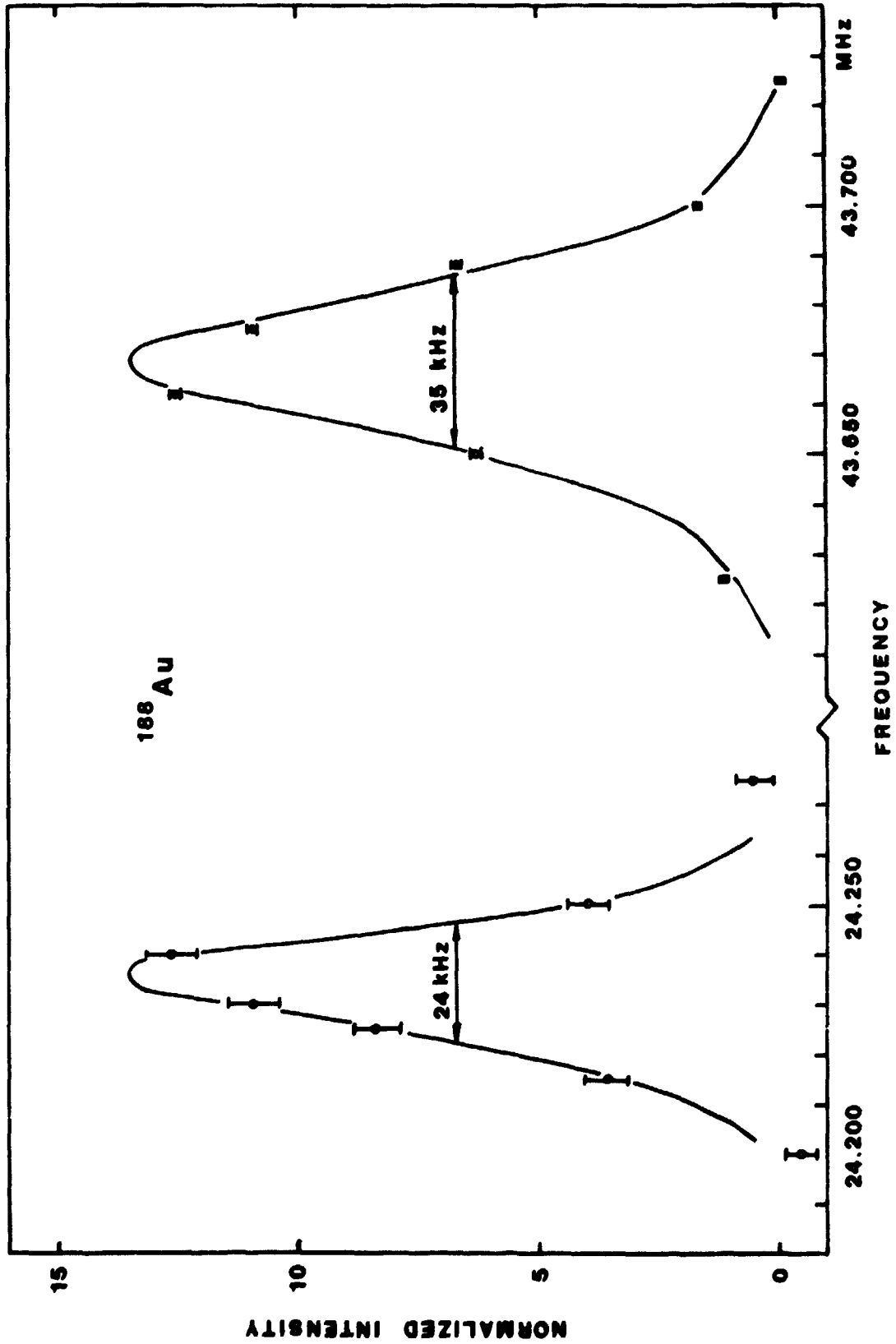


Fig. 2

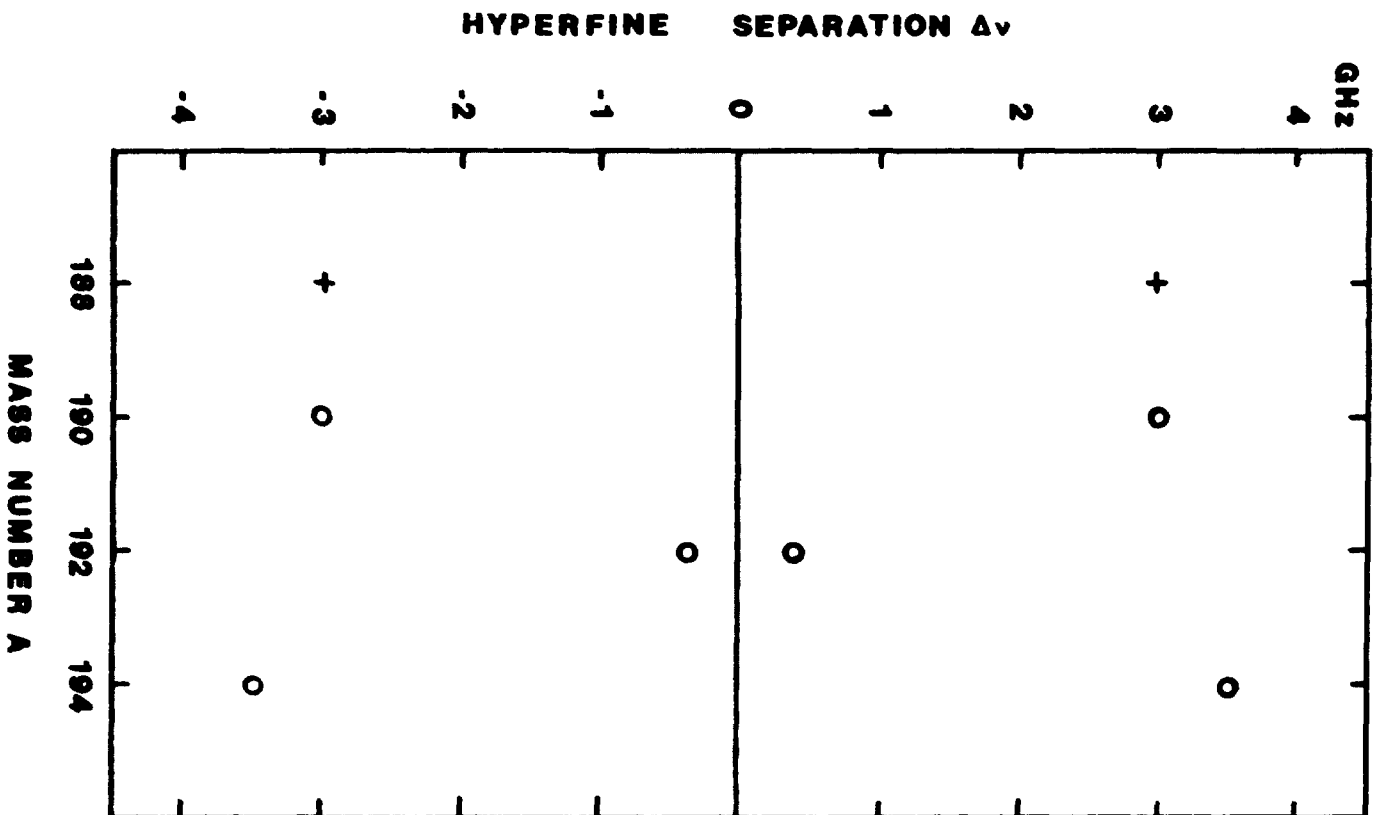


Fig. 3

