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FEW-BODY SYSTEM AND PARTICLE RESONANCES *

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ABSTRACT

Techniques of few-body system in nuclear physics are exploited to analyse the spectrum of the T resonance and its family. Their relation to nuclear resonances are established so as to apply few-body dynamical techniques in the dynamical structure of particles carrying the truth quantum number.

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After the recent discovery of T resonance ¹⁾ (upsilon meson) around 9.4 (GeV), it is now widely accepted that the T family is a bound state $t\bar{t}$ of a fifth quark, carrying the truth quantum number. In this note, we are concerned with some of the few-body aspects of this spectrum. The discovery of T resonances and the speculation of the fifth quark have given rise to a rich spectroscopy of elementary particles and has made $SU(5)$ symmetry topical recently. ²⁾ We have given a new spectroscopy for these resonances incorporating the truth quantum number with charge $\frac{2}{3}$. ³⁾ The mass spectra in GeV of this family of mesons and the corresponding new baryons carrying the truth quantum number that are identified in our scheme are as follows:

Mesons $T = 9.44$ (GeV), $T' = 9.69$ (GeV), $T'' = 9.86$ (GeV),

$T''' = 10.05$ (GeV), $T^0 = 10.37$ (GeV) .

Baryons $uut = 17.06$ (GeV), $ust = 17.20$ (GeV), $cut = 17.23$ (GeV),

$dct = 17.0$ (GeV), $ttt = 17.46$ (GeV) .

From the energy level diagram, it becomes evident that the level of the hadronic states increases as the content of the t quark increases, in the composite structure of the topomonium hadrons. This copes with the Eichten and Gottfried diagram ⁴⁾. The level diagrams of these truthful hadrons show remarkably the salient aspects of the typical nuclear physics level diagrams. The t quark is very heavy and the motion becomes non-relativistic. In this way this model contributes a great deal to our knowledge to relate such hadronic and nuclear physics spectroscopies.

The production of resonances and their properties in three and more particle systems appears to be a common feature in nuclear and particle physics. Resonances like the A_1 bump and spectators A seem to be the isobar content of the excited nuclear states. This then becomes pertinent with respect to the quark model description of hadrons. The quark model predicts two octets of vector mesons with $J^{PC} = 1^{++}$ and 1^{+-} , in which A and B mesons are accommodated. Further, the interaction of relativistic heavy ions has some of the following properties. The existence of many closely packed $J^{PC} = 1^{--}$ states around 9.4 GeV demands model states of excitations that are prevalent in nuclear states having the same spin and parity. The excitations of the great dipole resonance in light nuclei are described as the collective excitations

of many particle-like hole states $J^P = 1^-$, differing only in shell model orbitals of the various states. Analogously, the states of the T family can be imagined as five-quark composite, dynamically paired as $(t\bar{q})$, where q 's are $SU(4)$ quarks. These states represent the exotic mesons. The exotic quarks and antiquarks EE now become eight- or six-dimensional representations to cope with the data, and the E mesons have a type of string structure ⁵⁾. The colour and quark confinement string models must become phenomenological in such situations ⁶⁾ and resemble giant dipole resonant states. Hence the few-body techniques can also be aptly exploited for the dynamical calculations of these resonant mesons. The colour and quark confinement in a deformed MIT bag model for these resonances resemble the unified model of nuclear structure.

The hadronic decay of heavy quark Q and anti-quark \bar{Q} vector mesons proceeds via the annihilation of the $Q\bar{Q}$ system into gluons, which materialize into $SU(5)$ hadrons. The annihilation rate of $Q\bar{Q}$ into gluons calculated in the lowest-order quantum chromodynamic perturbation theory is an analogue of the Ore-Powell formula for orthopositronium decay, being studied in nuclear physics ⁷⁾. The three-body decay analogue, namely the three-gluon decay for these resonances via the three-gluon channel, becomes clear only in the case of hadrons containing heavy quarks such as T and ψ resonances, and this picture becomes clear for T particles. This may further improve if we proceed to hadrons containing still heavier quarks.

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