

A SYNCHRONOUS BEAM SWEEPER FOR HEAVY IONS

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The Argonne Tandem Linac Accelerator System (ATLAS) facility at Argonne National Laboratory provides a wide range of accelerated heavy ions from the periodic table. Frequently, the beam delivery rate of 12 MHz is too fast for the type of experiment on line. Reaction by-products from a target bombardment may have a decay interval much longer than the dead time between beam bunches. To prevent data from being corrupted by incoming ions a beam sweeper was developed which synchronously eliminates selected beam bunches to suit experimental needs. As the SWEEPER is broad band (DC to 6 MHz) beam delivery rates can be instantaneously changed. Ion beam bunches are selectively kicked out by an electrostatic dipole electrode pulsed to 2 kVDC. The system has been used for almost three years with several hundred hours of operating time logged to date. Beam bunch delivery rates of 6 MHz down to 25 kHz have been provided. Since this is a non-resonant system any beam delivery rate from 6 MHz down to zero can be set. In addition, burst modes have been used where beam is supplied in 12 MHz bursts and then shut down for a period of time set by the user.

Beam Pulse Removal Technique

The ATLAS accelerator system is synchronized from a master clock operating at 12-125 MHz. All accelerating structures are synchronized to this clock and therefore are harmonically related. Accelerated ions arrive at a rate of 12.125 MHz and typically have a bunch width of 1.0×10^{-9} second at the beam sweeper station. We use electrostatic fields to synchronously DEFLECT selected ion bunches onto a pair of vertical slits located a short distance downstream of the beam sweeper. The non-deflected ion bunches continue on through to be accelerated to higher energies. Voltage pulses on the deflection electrode are synchronized to the master clock, therefore rise and fall times must be less than one clock period. This is necessary to avoid unwanted vertical forces on the beam bunches selected for further acceleration. The sweeper must drive the deflection electrode to sufficient voltage for all species of ions and various charge states. The amount of ion deflection is a function of many factors as shown below:

$$\Delta Y = \frac{L Q_s V_p e}{S(Q_T + 1) V_T} \quad (1)$$

MASTER

where: ΔY = vertical displacement of beam bunch at
vertical slits (meters)

L = drift distance to vertical slits
(meters) 2.06 M

Q_s = charge state of ions within deflecting
electrode 5

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Sweeper Output Stage

Microwave planar triodes (8940) were selected to drive the deflection electrode because of their low parasitic capacitance, unique plate current characteristics, and modest grid drive requirements. Although these vacuum tubes have high peak plate current capability, their average grid dissipation is low. For this reason considerable effort was made to reduce stray capacitance in the system. The deflection electrode capacitance and associated stray capacitance due to the vacuum tubes, transformers, and mechanical mounting all add together. The total capacitive load must be driven to 1 kV or more, and then discharged in less than 100 nsec. This must be done at frequencies up to 4.0 MHz in our present configuration. The planar triodes control grid dissipation is limited to 2.0 watts average power. This limit is easily exceeded when operating at high switching frequencies, so load capacitance must be minimized. To charge a given load capacitance in a short interval of time requires a specific amount of vacuum tube plate current.

$$I = \frac{CAV}{T} \quad (2)$$

where: ΔV = voltage change across the load in volts
(1000V)

C = load capacitance in farads
(7.15×10^{-11} F)

T = charging time in seconds
(5×10^{-8} sec)

I = Vacuum tube plate current in amperes.

In our beam sweeper, charging the deflection electrode to 1000 V in 50 nsec required 1.43 amperes of plate current from tube V_1 . To discharge this energy vacuum tube V_2 must sink the same amount of current in 50 nsec or less. Considerable effort was made in reducing all capacitances that contribute to the total load.

The various load capacitances are shown below:

V_1 filament transformer capacitance = 6.5 pf

V_1 grid drive transformer capacitance = 6.5 pf

V_1 plate to cathode capacitance = 10.0 pf

V_1 mounting capacitance	= 10.5 pf
V_2 plate to cathode capacitance	= 10.0 pf
Beam deflection electrode capacitance	= <u>28.0 pf</u>
Total load capacitance	= 71.5 pf

Referring to Fig. 2, tube V_1 charges the deflection electrode to a specified level and tube V_2 discharges it. Vacuum tube V_2 must be held in a cut-off condition during V_1 conduction and remain cut-off until needed to discharge the electrode. The grid drive signal for tube V_2 is a composite of two signals. First, when the electrode must be discharged, a positive going grid drive pulse of short duration is used to drive V_2 into sufficient conduction. After that, V_2 grid bias is maintained at zero volts for as long as needed. Just before tube V_1 is turned on, the grid bias on V_2 is driven to a cut-off level. The composite grid drive signal for V_2 is shown in Fig. 2 along with other related signals.

Both grid drivers use field effect transistors because these devices are low on resistance and are relatively easy to drive. The vacuum tubes are used in a common cathode configuration which means that Miller effect must be dealt with. As you know, Miller effect is the apparent multiplication of triode tube input capacitance due to grid-to-plate coupling and voltage gain μ .

$$C_{IN} = C_{GK} + C_{GP} (1 + \mu) \quad (3)$$

where: C_{IN} = input capacitance modified by Miller effect

C_{GK} = grid to cathode capacitance

C_{GP} = grid to plate capacitance

μ = control grid voltage gain with respect to plate (with plate current constant).

For the tubes used, the Miller effect input capacitance is about 900 pf. The result of this phenomenon is that fast control grid voltage changes are difficult to achieve. The use of medium power n-channel enhancement mode field effect transistors as grid drivers overcomes the Miller effect by brute force.

Vacuum tube plate dissipation is a function of operating voltage capacitive load and switching frequency.

$$P = \frac{1}{2} CV^2f$$

where: P = plate dissipation of either vacuum tube
in watts

C = total load capacitance in farads

V = peak electrode voltage in volts

f = switching frequency in hertz.

At 4.0 MHz switching frequency each vacuum tube will dissipate:

$$P = \frac{1}{2}(71.5 \times 10^{-12} \text{F})(1000\text{V})^2(4 \times 10^6 \text{Hz}) = 143 \text{ watts (5)}$$

Switch tube V_1 has a string of zener diodes and by-pass capacitors in series with its' cathode lead. The zener diodes develop a relatively constant bias voltage to keep V_1 cut-off when V_2 is conducting. Grid drive for V_1 (Point C in Fig. 2) is of sufficient magnitude to overcome the zener diode bias and drive V_1 into saturation. The 50Ω series resistor in V_1 grid circuit serves to damp oscillations and limit grid dissipation.

The tubes are fitted with isolated water cooling jackets and fans provide air cooling to the tube sockets. Figure 3 shows a photograph of the deflection electrode signal.

Acknowledgements

This work is a continuation of a concept first developed by Kenneth W. Johnson and Benny E. Clift (Argonne).

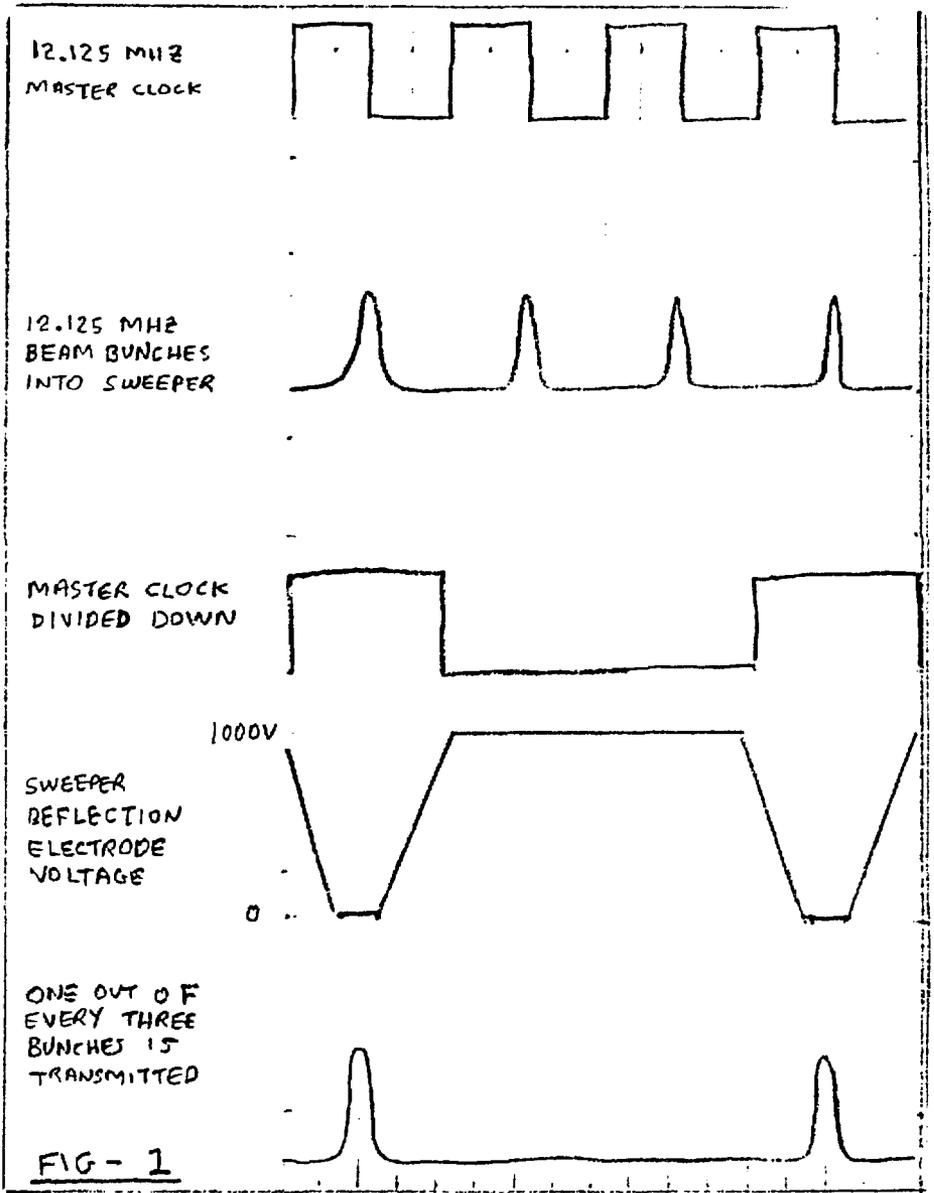
Work supported by the U. S. Department of Energy, Nuclear Physics Division, under Contract W-31-109-ENG-38.

Figure Captions

- Fig. 1 Sweeper timing relationship.
- Fig. 2 Simplified beam sweeper schematic.
- Fig. 3 Sweeper output voltage.

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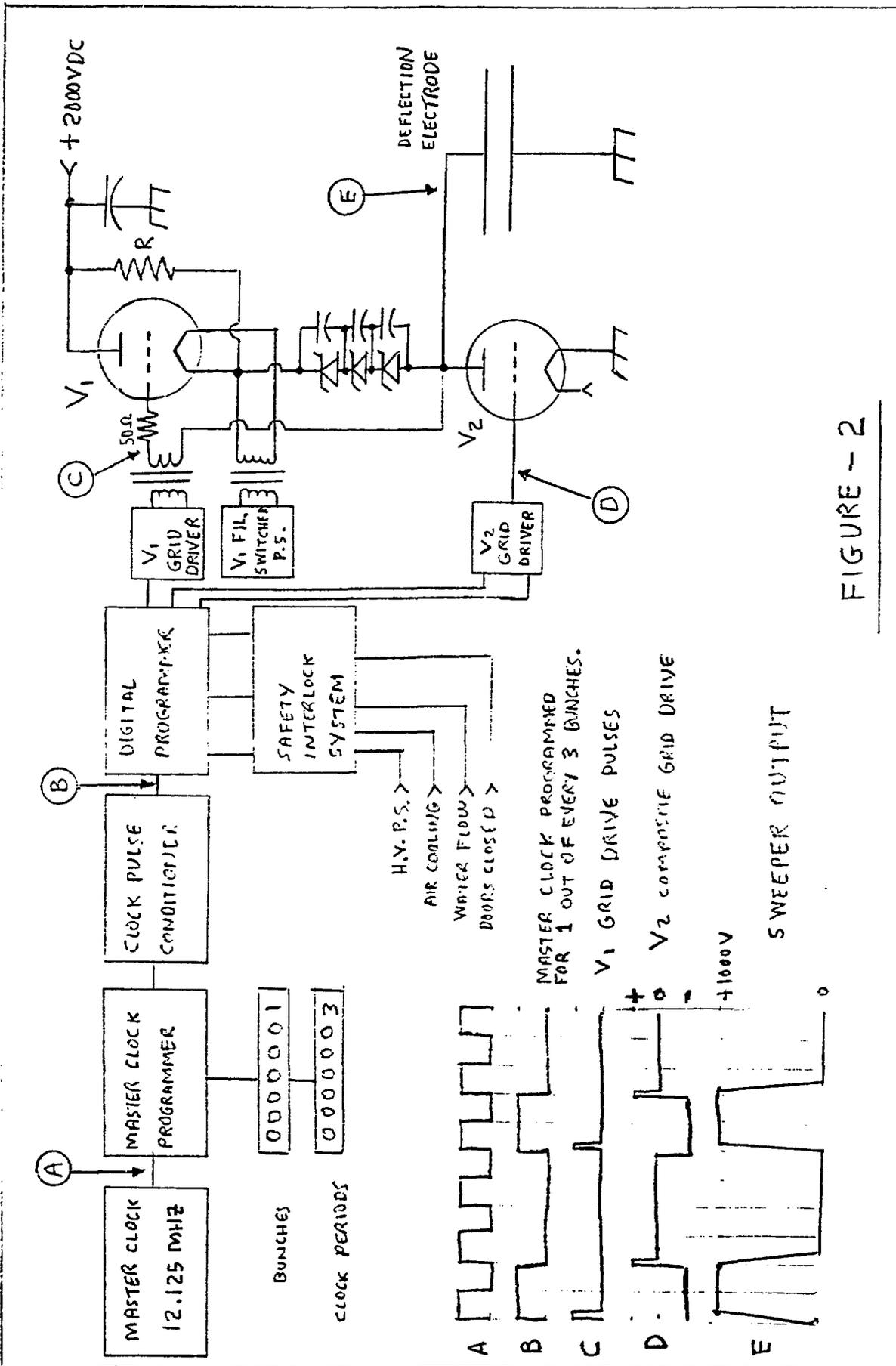
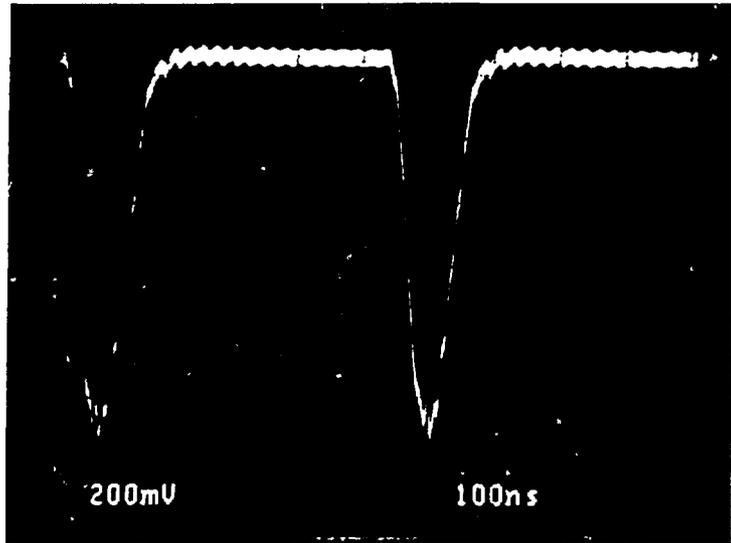


FIGURE - 2



SWEEDER OUTPUT VOLTAGE

200 V/cm

100 NS/cm

FIGURE-3