

High Ion Charge States in a High-Current, Short-Pulse, Vacuum Arc Ion Source

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Abstract

Ions of the cathode material are formed at vacuum arc cathode spots and extracted by a grid system. The ion charge states (typically 1-4) depend on the cathode material and only little on the discharge current as long as the current is low. Here we report on experiments with short pulses (several μs) and high currents (several kA); this regime of operation is thus approaching a more vacuum spark-like regime. Mean ion charge states of up to 6.2 for tungsten and 3.7 for titanium have been measured, with the corresponding maximum charge states of up to 8+ and 6+, respectively. The results are discussed in terms of Saha calculations and freezing of the charge state distribution.

I. Introduction

Broad beams of multiply charged metal ions can be produced by vacuum arc plasma sources. They have been described in detail in [1, 2], and we give here only a brief description. An arc between metal electrodes in vacuum is usually initiated by applying a positive high voltage pulse to a trigger electrode. The arc burns in the plasma formed at non-stationary cathode spots. The cathodic plasma expands from the cathode spots into the vacuum and reaches the ion extractor system which is usually a three-grid system of the acceleration-deceleration type. It is common that ion sources of this type operate in a repetitively pulsed mode (arc duration of order 1 ms or less) with repetition rates of a few pulses per second, an arc current of 100-300 A, an ion beam current of order 100 mA (during pulses), and an extractor voltage in the range 20-100 kV. The ion charge state distribution (CSD) was investigated in a number of papers [1-4], and the most recent table is given in Ref. [1] for 50 cathode materials.

CSDs were found to be fairly reproducible in different embodiments of the vacuum arc ion source. However, it turned

out that the CSDs do not only depend on the cathode material but also on the time elapsed since the arc pulse was started [5, 6], and on the presence and strength of an external magnetic field [6-9].

In the present paper we investigate the influence of the discharge current on the CSD. This is motivated by the fact that very high ion charge states have been observed spectroscopically in short-pulse, high-current vacuum spark experiments [10, 11]. High charge states are of interest for obtaining high ion energies at a given extractor voltage.

II. Experiments

For our experiments we used a vacuum arc ion source similar to what has been described previously [1, 12]. The arc was fed by a 7-stage LC-pulse line of 1.5 Ω impedance, resulting in pulses of 100-300 A amplitude, 5 μs rise time and 250 μs duration. The cathode spots burned on the front face of a water-cooled, cylindrical cathode of 6.2 mm diameter; the annular anode had an inner diameter of 20 mm and was 1 cm apart from the cathode surface (Fig. 1). The anode hole diameter was

reduced to 4 mm or 2 mm for discharge currents higher than 300 A and 3 kA, respectively, in order to avoid breakdown of the extractor voltage. The plasma leaving the arc discharge region expanded while streaming towards the extractor grids. The extractor grids had holes of 6 mm diameter and were placed 10 cm downstream the anode. The ion beam had a diameter of about 10 cm.

To achieve higher currents and shorter pulses, the LC pulse forming network was replaced by a fast high voltage capacitor ($C = 12 \mu\text{F}$ or $100 \mu\text{F}$ in different experiments). The arc current was controlled by using various inductive loads ($0.1\text{-}10 \mu\text{H}$) and a $100 \text{ m}\Omega$ Ohmic resistor. In most

experiments, the capacitor was located very close to the electrode arrangement to keep the circuit inductance relatively small (about $0.1 \mu\text{H}$), and thus pulsed currents of several kiloamperes could be obtained even with the charging voltage not exceeding 600 V. The arc pulse length was in the range of $3\text{-}25 \mu\text{s}$ (FWHM). Ions of charge state z are accelerated to the energy $E_z = z U_{extr}$; the default extractor voltage was $U_{extr} = 50 \text{ kV}$. Titanium and tungsten were used as cathode materials. The CSD was investigated by a time-of-flight (TOF) charge-to-mass spectrometer described elsewhere [12].

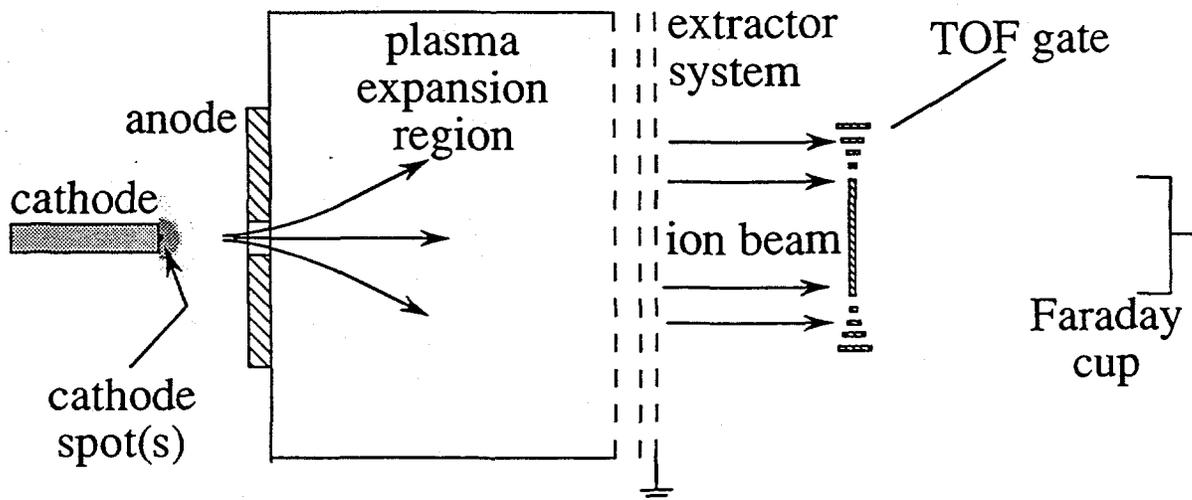


Fig. 1 Schematic of the source of multiply stripped metal ions and TOF diagnostics; (not to scale).

III. Results

Fig. 2(a) shows a typical TOF spectrum for Ti obtained by the "standard" low-current vacuum arc ion source. The arc current was 150 A, and the mean ion charge state was $\bar{z} = 2.2$, measured 100 μs after triggering. As a first step towards higher currents and shorter pulses we used a $100 \mu\text{F}$ capacitor and a circuit inductance of $3.6 \mu\text{H}$. The titanium CSD at a peak current of 560 A (30 μs after triggering) is shown in Fig. 2(b);

$\bar{z} = 2.4$ that is only slightly higher than at "standard" operation.

Higher discharge currents were achieved by reducing the inductive load. Figure 2(c) shows the titanium CSD at 1.8 kA peak current ($12 \mu\text{F}$ capacitor, the FWHM of the current pulse was 6 μs) with $\bar{z} = 3.6$, i.e. significantly higher than at "standard" operation, and $\bar{z} = 3.75$ has been obtained at 3.4 kA.

It is known from previous work [3] that cathode materials of high boiling point tend to

have high charge states; the empirical law $\bar{z} = 0.38 (T_{bp}(\text{°C})/1000) + 1$ was derived for discharge currents smaller than 200 A. To proof this rule at higher currents we replaced the titanium cathode ($T_{bp} = 3287 \text{ °C}$) by a tungsten cathode ($T_{bp} = 5660 \text{ °C}$). The mean ion charge state at "standard" operation (150 A, 100 μs after triggering) is $\bar{z} = 3.1$. Figs. 3(a) and (b) show TOF spectra of tungsten ion beams obtained with 12 μs pulses at a peak current of 3.4 kA with $\bar{z} = 5.8$ and 6.3. There is obviously a considerable shot-to-shot variation in the CSD but a clear increase of \bar{z} with current (Fig. 4).

IV. Discussion

The formation of ions at cathode spots is associated with phase transitions and ionization of the cathode material. Although a number of papers have been published [13-19], a satisfying model has not yet been formulated. It has been shown that inhomogeneous heating of the cathode surface (mainly by ion bombardment) leads to local melting and the formation of a dense, non-ideal plasma of micron-size. This is the plasma of "microspots" which expands rapidly, driven mainly by the pressure gradient to the vacuum ambient. While local thermodynamic equilibrium is established close to the spot (within a distance of order 100 μm from the spot center), the expanding plasma quickly goes into non-equilibrium at low density.

As long as equilibrium conditions are fulfilled (the time between inelastic collisions is short compared to the characteristic expansion time), the plasma composition can be calculated by a set of Saha equations [15, 17, 19, 20]. Figs. 5-8 show calculated plasma compositions assuming a plasma temperature of 3 eV and 5 eV for titanium and 4 eV and 6 eV for tungsten. For these calculations, ionization energies and partition functions have been taken from Refs. [21-23]. For the lowering of the ionization energies we used the Debye-Hückel model which is valid for weakly non-ideal plasmas.

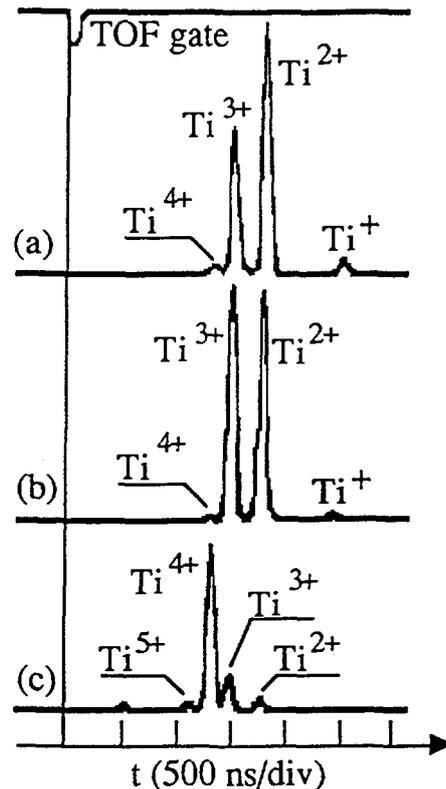


Fig. 2 TOF spectra of titanium ion beams: (a) "standard" ion source, 150 A, 100 μs after triggering; (b) 30 μs after triggering, using a 100 μF capacitor, 560 A, 60 μs pulse; (c) 3 μs after triggering, 12 μF capacitor, 1.8 kA; 6 μs pulse.

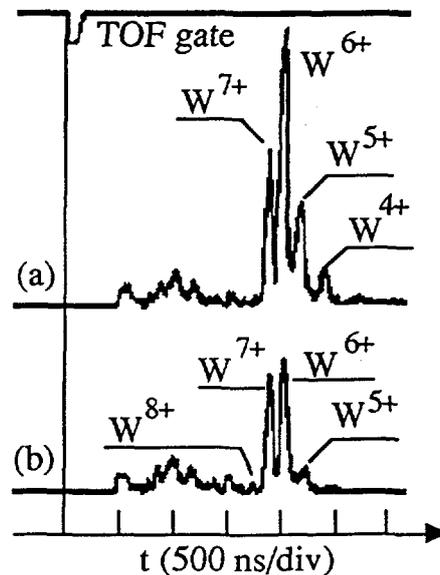


Fig. 3 TOF spectra of tungsten ion beams: (a) 6 μs after triggering, 100 μF capacitor, 3.4 kA; 12 μs pulse; (b) same as (a) showing the large shot-by-shot variation.

The calculation for tungsten are somewhat uncertain since we used self-consistent-field Hartree-Fock results for the ionization energies [21] and estimates for the partition functions (to our knowledge, no experimental data are available).

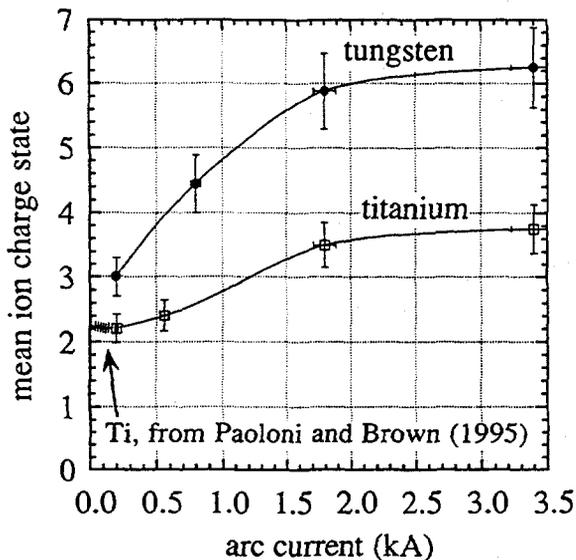


Fig. 4 Increase of the mean ion charge state with discharge current. Each point represents an average of 16 discharges.

When ionizing/recombining collisions become infrequent the composition of the plasma “freezes”. Although “freezing” is actually a step-wise process [15], it is reasonable to look for a combination of plasma density and temperature which approximately gives the experimentally observed CSD. In this way, we determined a “freezing temperature” and “freezing density” as indicated in Figs. 5-8.

At low current, plasma is produced at a single cathode spot or very few spots. The plasma of each spot can expand independently from the plasma of other spots, and the plasma density decreases as

r^{-2} where r is the distance from the spot center. Freezing typically occurs at a density of heavy particles $n_h \approx 10^{25} \text{ m}^{-3}$ ([17] and Figs. 5, 6).

At high current, the freezing process and the resulting CSDs are changed by two effects. First, more energy is invested in each heavy particle. If we assume that (i) the amount of plasma produced is proportional to the current and (ii) the plasma conductivity does not change significantly with increasing current, then the energy per particle is approximately proportional to the current. Therefore, the plasma temperature increases with current. Second, many spots are simultaneously active at high current, and the plasma of each spot expands under the influence of neighboring plasmas. Thus the expansion is not as rapid as in the previous case, and the density drops as $r^{-\nu}$ with $\nu < 2$. The characteristic expansion time becomes longer, and freezing occurs at lower densities and larger distances from the cathode. These effects are illustrated in Figs. 7 and 8.

New effects must be expected at much higher currents (many kA). Erosion of material occurs not only from the cathode but also from the anode. Compression and plasma heating by the self-magnetic field leads to much higher charge states in “hot spots” within the plasma [10, 12]. More effort will be needed these effects for the formation of beams of even higher charged metal ions. A broad beam metal ion implanter based of multiply charged ions is under construction at LBNL.

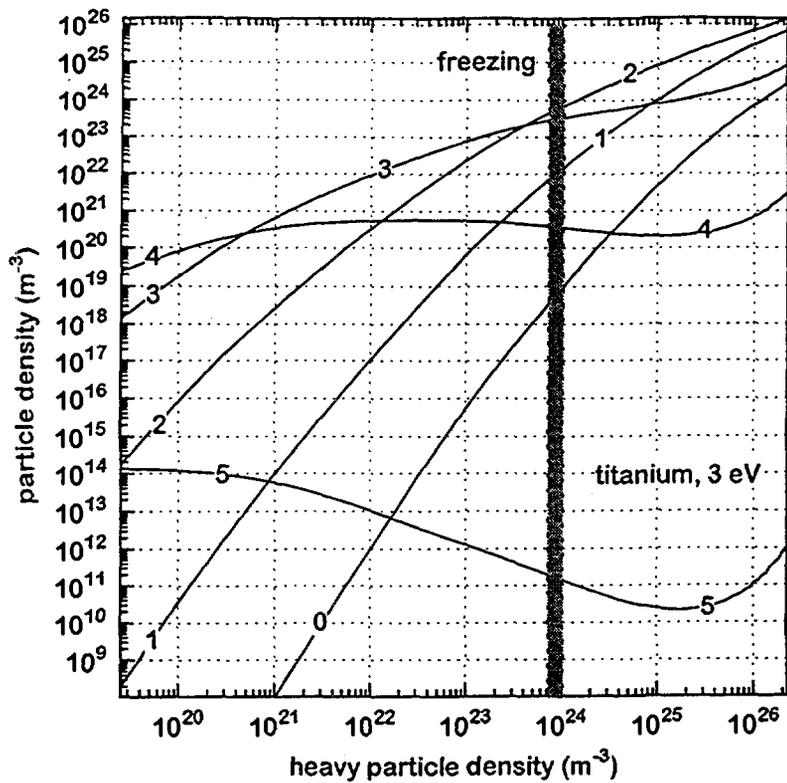


Fig. 5 Composition of a Ti plasma in equilibrium at 3 eV (Saha calculation); the freezing in the case of low current is indicated.

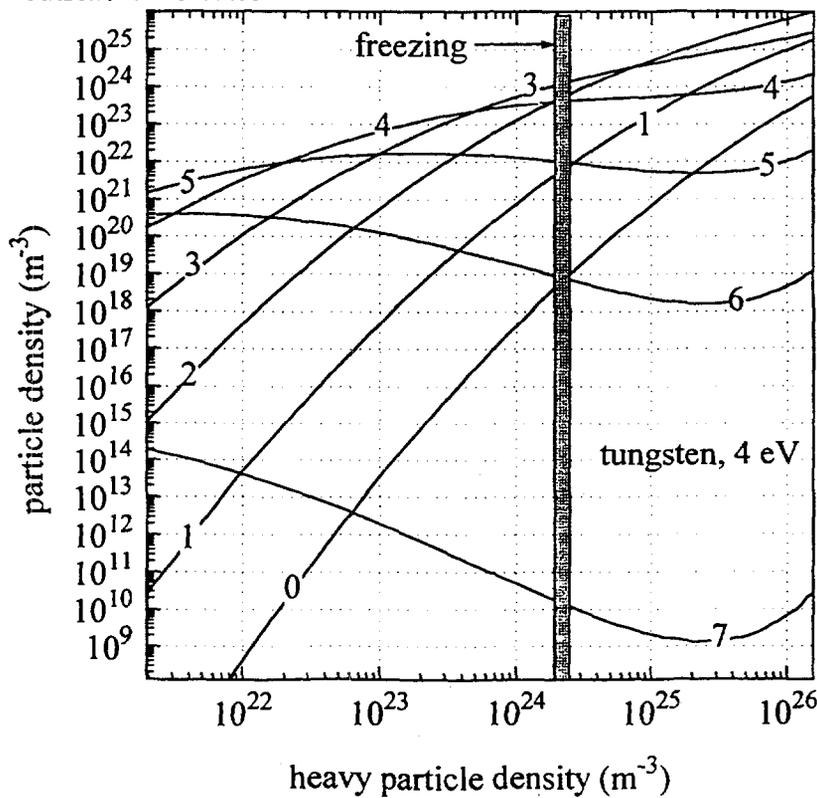


Fig. 6 As Fig. 5 but for tungsten at 4 eV.

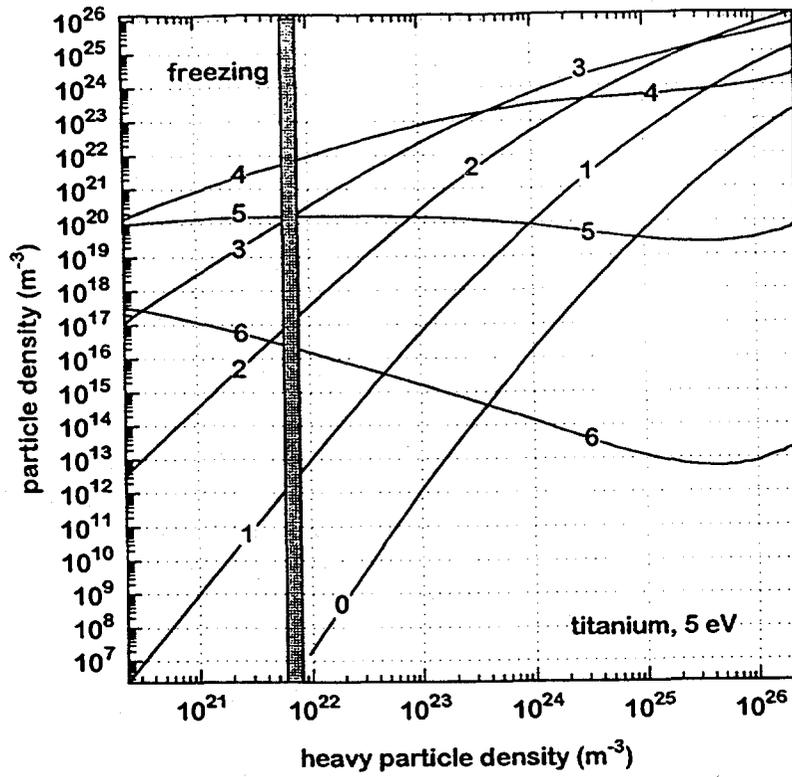


Fig. 7 Composition of a Ti plasma in equilibrium at 5 eV; the freezing in the case of high current is indicated.

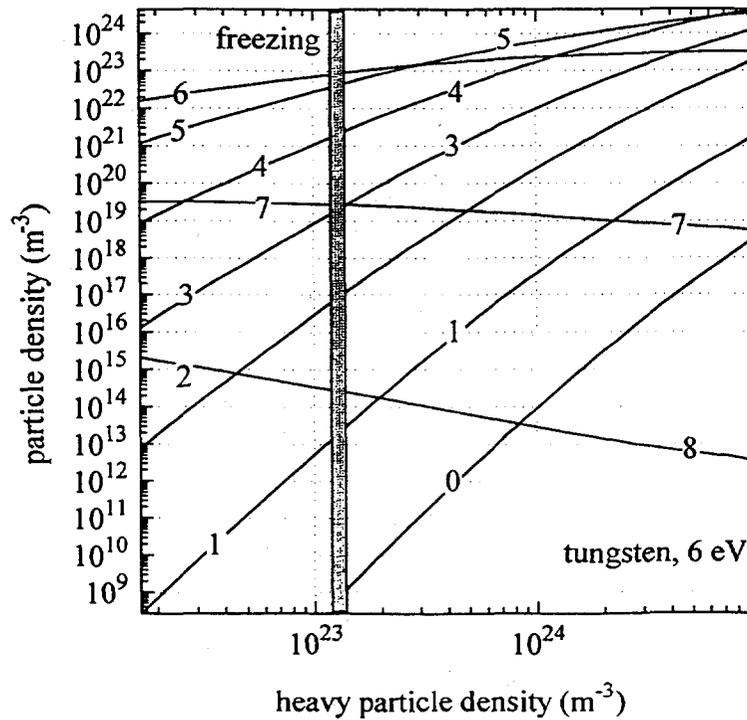


Fig. 8 As Fig. 7 but for tungsten at 6 eV

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