

## JET Radiative Mantle Experiments in ELMy H-Mode

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**Abstract:**

Radiative mantle experiments were performed on JET ELMy H-mode plasmas. The Septum configuration was used where the X-point is embedded into the top of the Septum. Argon radiated 50% of the input power from the bulk plasma while  $Z_{\text{eff}}$  rose from an intrinsic level of 1.5 to about 1.7 due to the injected Argon. The total energy content and global energy confinement time decreased 15% when the impurities were introduced. In contrast, the effective thermal diffusivity in the core confinement region ( $r/a = .4 - .8$ ) decreased by 30%. Usually, JET ELMy H-mode plasmas have confinement that is correlated to the edge pedestal pressure. The radiation lowered the edge pedestal and consequently lowered the global confinement. Thus the confinement was changed by a competition between the edge pedestal reduction lowering the confinement and the weaker RI effect upon the core transport coefficients raising the confinement. The ELM frequency increased from 10 Hz Type I ELMs, to 200 Hz type III ELMs. The energy lost by each ELM reduced to 0.5% of the plasma energy content.

Highly radiating plasmas occur when high Z impurities such as Neon, Argon, Nitrogen, and Krypton are injected into tokamak plasmas. The consequential reduction of power flowing to the first wall components has potential for improving reactor power handling [1]. However, experiments feature different consequences to the plasma energy confinement [2]. On the one hand, impurity injection into TEXTOR [3], ISX -B [4] and DIII-D [5] improved confinement, while on the other hand, impurity injection into JET [6] degraded confinement. These experiments feature many differences, (such as the use of limiters on TEXTOR and ISX-B and divertors on the other machines) which could be playing some role in the different experimental results obtained.

This paper reports results from JET impurity seeding in ELMy H-mode plasmas. Our philosophy was to have a joint experimental program including many experimentalists with experience of radiative mantle experiments. Our strategy was to use their experience to improve the JET ELMy H-Mode, since that regime is a reactor relevant regime that could benefit from reduced power flows to the divertor. Our intention was to perform the experiment in a similar manner as TEXTOR (albeit in a different plasma regime and with a different plasma boundary) and thus assess the origin of the different results.

The experiment was limited to a two-day duration, and so can not be considered to be optimised. However, plasmas were obtained which followed the techniques used on TEXTOR and DIII-D. We observed:

- (i) Confinement degradation similar to that seen previously on JET [6], *however higher densities were achieved*
- (ii) The confinement degradation correlated with the reduction in the edge pedestal pressure as seen commonly on JET [7]
- (iii) The core ( $r/a = .4 - .8$ ) confinement improved with the impurity seeding similar to TEXTOR [3]
- (iv) The radiative fraction was high, and the energy lost per ELM was substantially reduced
- (v) Even during the impurity seeding, the  $Z_{\text{eff}} < 2$ .

28 discharges (ELMy H-Modes at 5 MA and 2.5 T heated by 12 MW neutral beams) with Argon or Neon radiating mantles were performed. The Argon produced

about twice as much core radiation as the Neon. Two divertor configurations were tried: 1) The Corner configuration (with the strike points at the cryo-pump port location) should maximise the divertor's impurity pumping capability, and 2) The Septum configuration with the X-point embedded into the Septum top. This configuration is like a pump limiter, since the plasma is bounded by a material surface but the SOL has access to the divertor pumps. JET Septum operation features an H-mode power threshold 2-3 times lower than other divertor configurations. Thus Septum operation has the conceptual (but not yet proven) advantage of allowing more radiation without reducing the power flow through the X-point to below the H-mode power threshold.

The bulk of the experimental campaign was devoted to modification of the deuterium and Argon gas injection rates in Septum plasmas. The best results came with  $3 \cdot 10^{22}$  electrons/s from  $D_2$  gas puffing, and a step-down evolution of Argon gas (Fig. 1 and 2). The radiating plasma is in quasi-equilibrium with stationary parameters lasting about 5 energy confinement times (Fig. 2). The difference between the present campaign, and previous JET highly radiating plasma experiments [6] is primarily that higher deuterium gas fuelling was used in conjunction with the Argon puffing (Fig. 1), and that the Septum configuration was explored.

Both TEXTOR [8] and DIII-D publications describe two additional factors important for confinement enhancement during radiative mantle experiments, namely:

1. Avoiding MARFE formation, and
2. The plasma edge distance to the outer wall.

In these JET experiments, variations of the position of the outer edge did not affect confinement and MARFEs did not occur in any of these plasmas.

The input power was about 12 MW from 17 to 23 sec in Fig. 1, and the radiated power rose approximately from 1MW bulk and 2MW X-point radiation to 5MW bulk and 3 MW X-point radiation about 1.2 sec after the Argon puffing began. The X-point radiation stayed about constant although less spiky as the ELM behaviour changed in the presence of the higher mantle radiation (Fig. 2). The ELMs had been regular type I ELMs of about 10 Hz frequency without Argon and became type III ELMs (200 Hz frequency). Due to the high density,  $Z_{\text{eff}}$  rose only modestly from about 1.5 without Argon to 1.7 with the Argon (Fig. 2). The central density rose about 10% to about  $8.5 \cdot 10^{19} \text{ cm}^{-3}$  with the Argon but the central electron temperature

fell from about 3 keV to about 2.5 keV. With the lower temperature, the  $d(d,n)^3\text{He}$  neutron emission fell from 3.5 to 2.5  $10^{15}/\text{sec}$ . The global (magnetic) confinement time fell 15%, although the TRANSP deduced effective conductivity coefficient at  $R=3.5$  m fell about 30% (Fig. 1).

The bolometer signals [9] were Abel unfolded to determine the radiation profile (Fig. 3) of the top half of the poloidal cross-section. The X-point radiation dominates the emission in the bottom half and obscures the core radiation pattern. The Argon impurity enhanced the radiation across the entire profile. The JET Argon bulk radiation profiles were similar to TEXTOR Neon profiles [8] although the JET radiation power density is about half that in TEXTOR. Both plasmas were at 85% of the Greenwald density limit, with similar beam heating power densities. The width of the JET edge radiation mantle is about half the TEXTOR Neon mantle width, or  $\frac{1}{4}$  the TEXTOR Argon width.

The fuel depletion by the seeded impurity was acceptably small ( $n_d/n_e = .75$ ) for JET as for TEXTOR (Fig. 2). Since the JET plasmas had large deuterium fuelling, the higher density caused high radiation levels at modest impurity injection rates (Fig. 1). By contrast, TEXTOR also had high density, but usually turn off the gas fuelling during the RI regime. Charge exchange spectroscopic observation of Carbon densities indicated the intrinsic impurity levels were not affected by the radiation mantle, and remained constant.

The Argon radiation mantle discharges achieved good confinement when compared with other JET plasmas taken at the same plasma conditions (Fig. 4). The comparison data in figure 4 span the earlier JET divertor configurations, comprising all the data in the official JET steady-state data bases having the quoted machine settings. The densities are 20% higher than have been otherwise achieved. Apparently, the radiative mantle was useful to achieve higher density as on TEXTOR and DIII-D but at the expense of a degraded confinement unlike TEXTOR and DIII-D. The reference plasma (no Argon, Septum configuration) also had slightly higher density than the Mark II plasmas (at these conditions). Possibly, the density is higher in part due to Septum operation as well as due to the radiation mantle.

The energy balance was analysed using the TRANSP code. The diagnostic consistency (comparison of the kinetic and magnetic energy contents) was not particularly good, with the kinetic energy exceeding the magnetic energy content by about 25%, for both the no Argon reference case and the selected Argon radiative mantle plasma. The coupling between the electrons and the ions was large due to the low electron temperatures and the high densities, so that separate electron and ion energy balances could not be distinguished. An effective  $\chi$  (total power loss described by a transport coefficient governed by the average of the electron and ion temperature gradient) was reduced by about 30% (Fig. 1 and 5) which is a similar reduction as was reported on TEXTOR [3] in this region of the plasma (Fig.5). The reduction exists in the calculated energy balance primarily because of the reduced conducted power flow through the core region due to the radiated power in figure 3. The previously mentioned difficulties inverting the radiation profiles also make the conduction coefficient uncertain in a magnitude that has yet to be determined.

The magnitude of the effective  $\chi$  is about  $0.5 \text{ m}^2/\text{sec}$  in the core confinement region (Fig. 5). The analysis of the time evolution of the soft X-ray emission indicates similar impurity particle transport in that spatial location (Fig. 5). However, the analysis of the impurity transport is uncertain since the Argon densities inferred from Charge Exchange recombination spectroscopy differ from those inferred from x-ray emission.

With the impurity radiation, the ELMs became more frequent and smaller (Figs. 2 and 6). The energy drop per ELM was about 0.5% of the plasma energy. This magnitude in a reactor would avoid significant erosion by instantaneous ELM energy deposition. Also, core radiation and the ELMs were losing nearly all the power from the plasma (Fig. 6), so that the power handling in the divertor is similar to detachment between ELMs [6] since little power is conducted directly to the strike point.

The no Argon reference plasma was at the high density, low edge electron temperature limit of the type I ELM behaviour [7]. The Argon radiation further lowered the edge electron temperature and slightly increased the edge density. The plasma electron pressure was approximately reduced about 30%. JET typically observes that the global energy confinement is related to the edge pressure [7], and these radiation mantle experiments follow that trend.



Thus the loss of confinement can be superficially understood as the typical JET energy confinement response to the reduction in the edge pressure [7]. The radiation mantle increased the ELM frequency causing more power to be lost out of the plasma volume by the ELM. This behaviour pattern has been quantified in several JET publications by characterising the ELM effect upon confinement [10,11]. However, for the case of the radiation mantles, the loss of confinement is not as large as was described by Fishpool [10], where a 200 Hz Elm frequency would have been accompanied by an H89 reduction to 1.4 instead of the observed 1.6 here. Apparently, the positive RI effect on the core confinement (Fig. 5) mitigates the JET empirical ELM or edge pedestal pressure reduction in confinement, but did not eliminate it completely.

A preliminary JET experiment used Argon radiation mantles to simulate the enhanced performance regimes found on ISX, DIII-D, and TEXTOR. Radiation mantles of about 50 % core radiation and 70 % total radiation were achieved. The experiments yielded three interesting results:

1. Radiative mantle plasmas were obtained with higher density (normalised to the Greenwald limit) when compared to JET plasmas at the same triangularity, field, and current.
2. The no-Argon, Septum configuration at high deuterium gas fuelling was the highest density, good confinement plasma in the JET data base at these conditions.
3. The role of the edge pressure in reducing the global confinement seemed to obscure the radiation induced improvement in the core confinement.

Correlation of the JET transport coefficients with shear stabilised ITG turbulence and gyro-Bohm scaling has not yet been studied. Possible future experiments include:

1. Further optimisation of the impurity type, deuterium and impurity fuelling rates, and divertor configurations is warranted.

2. Scaling of the pedestal and core transport coefficients might identify parameters where the core confinement effect can exceed the edge pedestal effect on global confinement.
3. Other regimes (besides the ELMy H- Mode) should be evaluated since they might become competitive with the ELMy H-Mode as a reactor regime.

## References

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## Figure Captions:

Figure 1. The time evolution of the Argon radiative mantle plasma 46980 (solid line) and the no-Argon reference plasma 46978 (dashed line). a) The input power was 12 MW, the volume radiation from the reference plasma was 1 MW, the Argon caused the volume radiation to rise to 5MW, while the X-point radiation remained at 2 MW. b) The total electron injection rate from deuterium (the reference plasma had slightly less deuterium fuelling) and from Argon. c) H $\gamma$ . d) The TRANSP calculated effective conductivity coefficient at R=3.5 m.

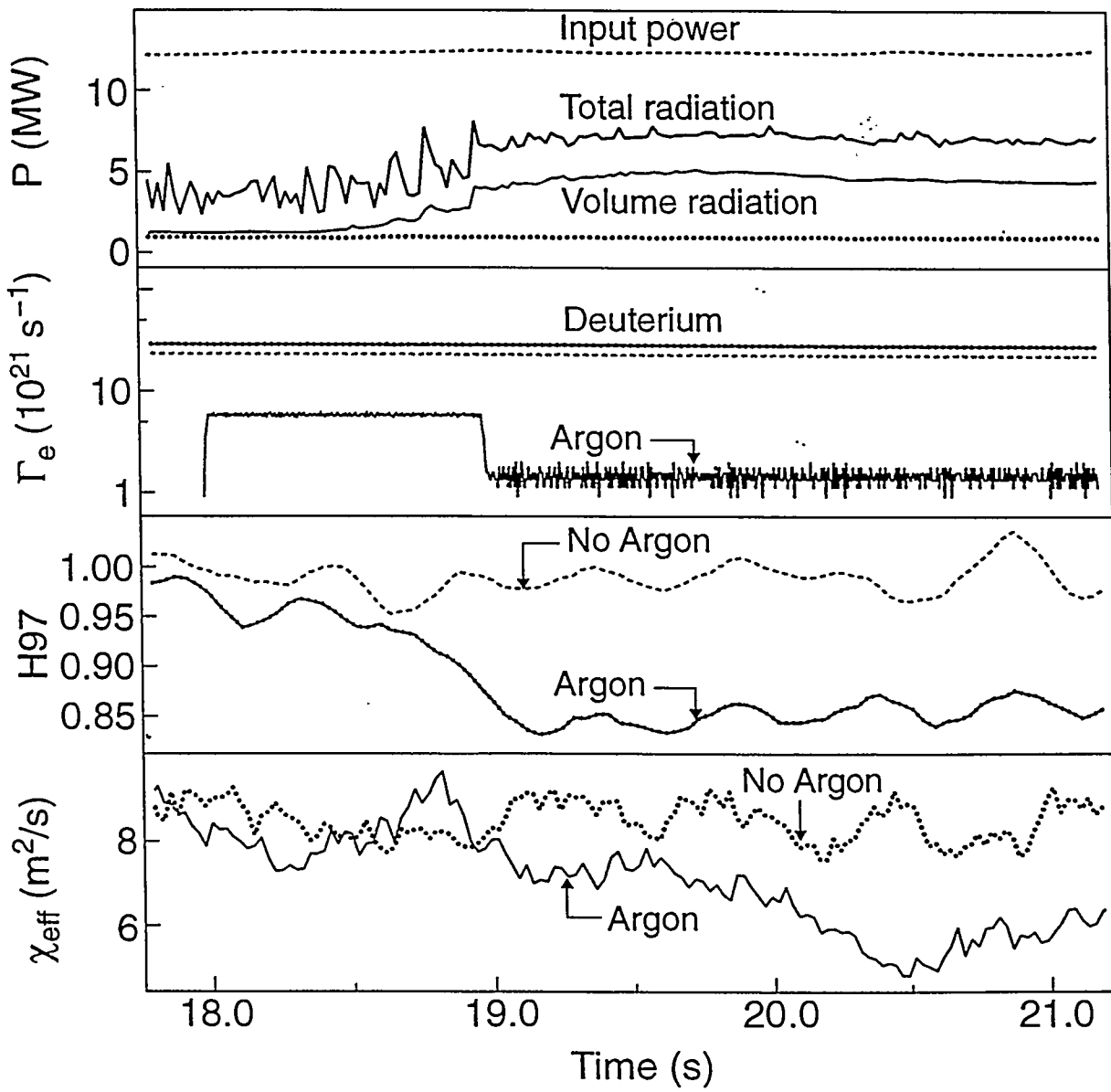
Figure 2. The time evolution of the Argon (46980) and reference plasma (46978) showing: a) the electron injection rate due to the Argon gas fuelling, b) the average  $Z_{\text{eff}}$  measured by visible Bremstrahlung, c) the Thomson scattering deduced central density, d) the ECE measured central electron temperature, and e) the H $\alpha$  light.

Figure 3. The electron temperature, density and Abel inverted radiation profiles overlaid with the TEXTOR profiles. Both TEXTOR and JET plasmas were near (.85 - .95) the Greenwald density limit, and heated by neutral beams to about the same heating power density. Neon or Argon radiating mantles [8] have the same temperature and density profile on TEXTOR. The radius is plotted as real distance from the plasma edge to compare the JET and TEXTOR radiating mantle widths. The JET radiation profile excludes the 3 MW of X-Point radiation by using only bolometer chords from the upper half of the plasma.

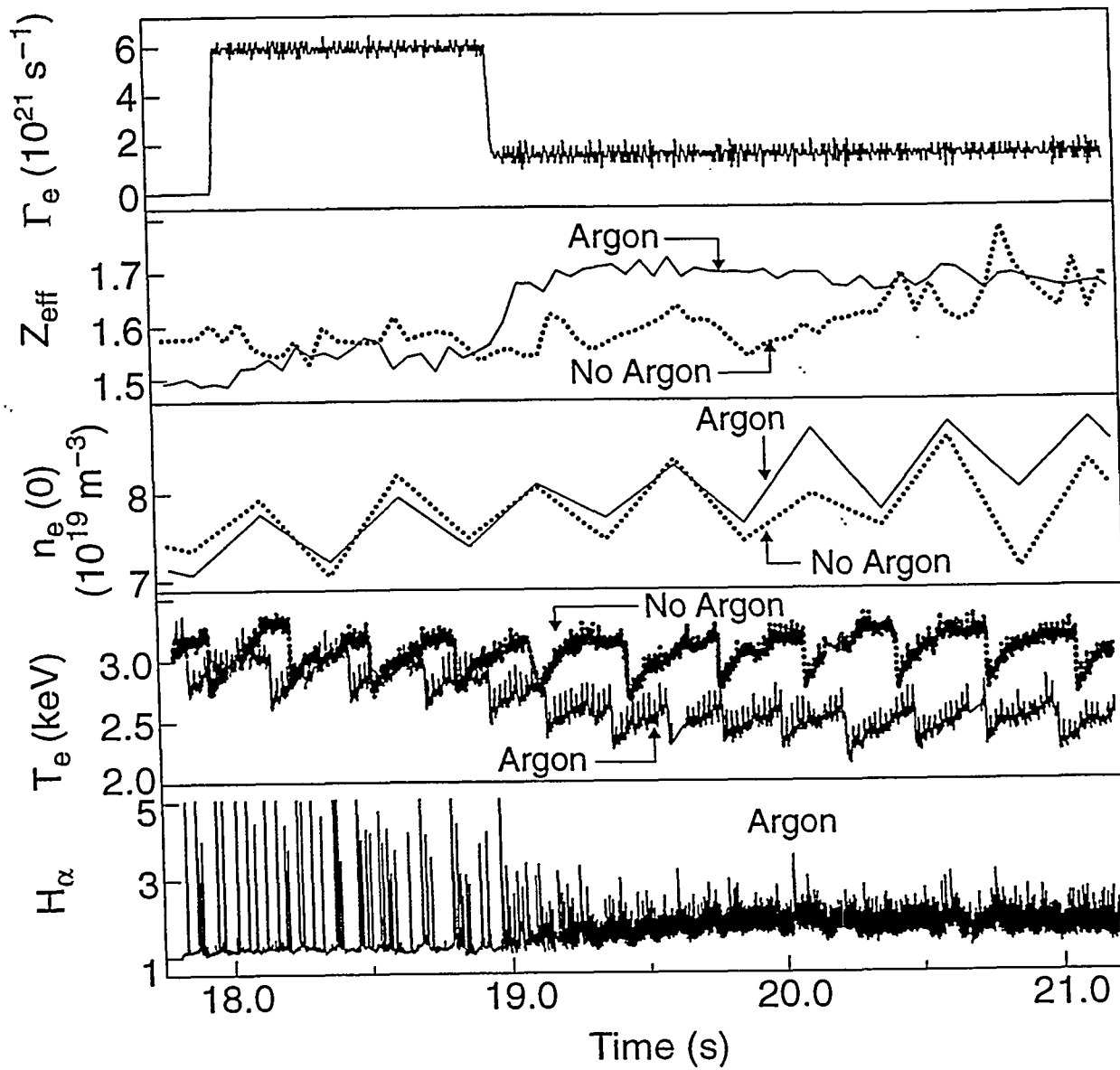
Figure 4. The normalised confinement plotted against the normalised density. The hollow data points are non-radiating JET plasmas at similar current, field, heating power, and triangularity. The Argon radiation mantle and reference plasmas (solid points) are the only ones in the Septum configuration. The solid curve is the JET empirical confinement degradation observed with density increases due to deuterium gas fuelling [7].

Figure 5. The profile of several transport coefficients indicates similar core confinement improvement in TEXTOR and JET. The TEXTOR electron chi is from Reference [3].

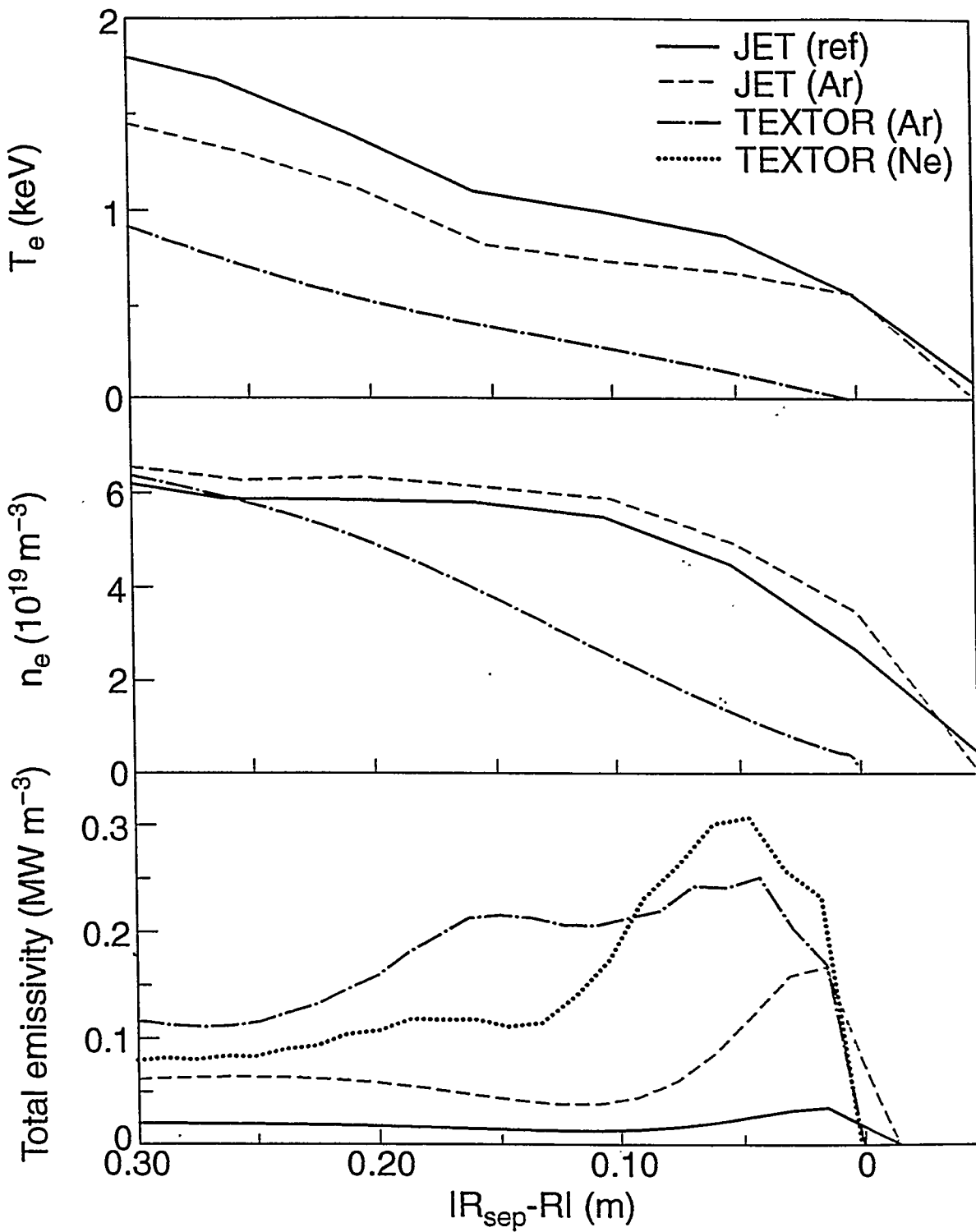
Figure 6. The fraction of total plasma energy lost at each ELM plotted as a function the observed ELM frequency. The dashed line is the amount of energy needed to be lost at each ELM for the ELMs to transport all the input power. The solid line is the amount of energy needed for the ELMs to transport  $\frac{1}{2}$  of the input power (since approximately  $\frac{1}{2}$  the input power is radiated in the plasma core).



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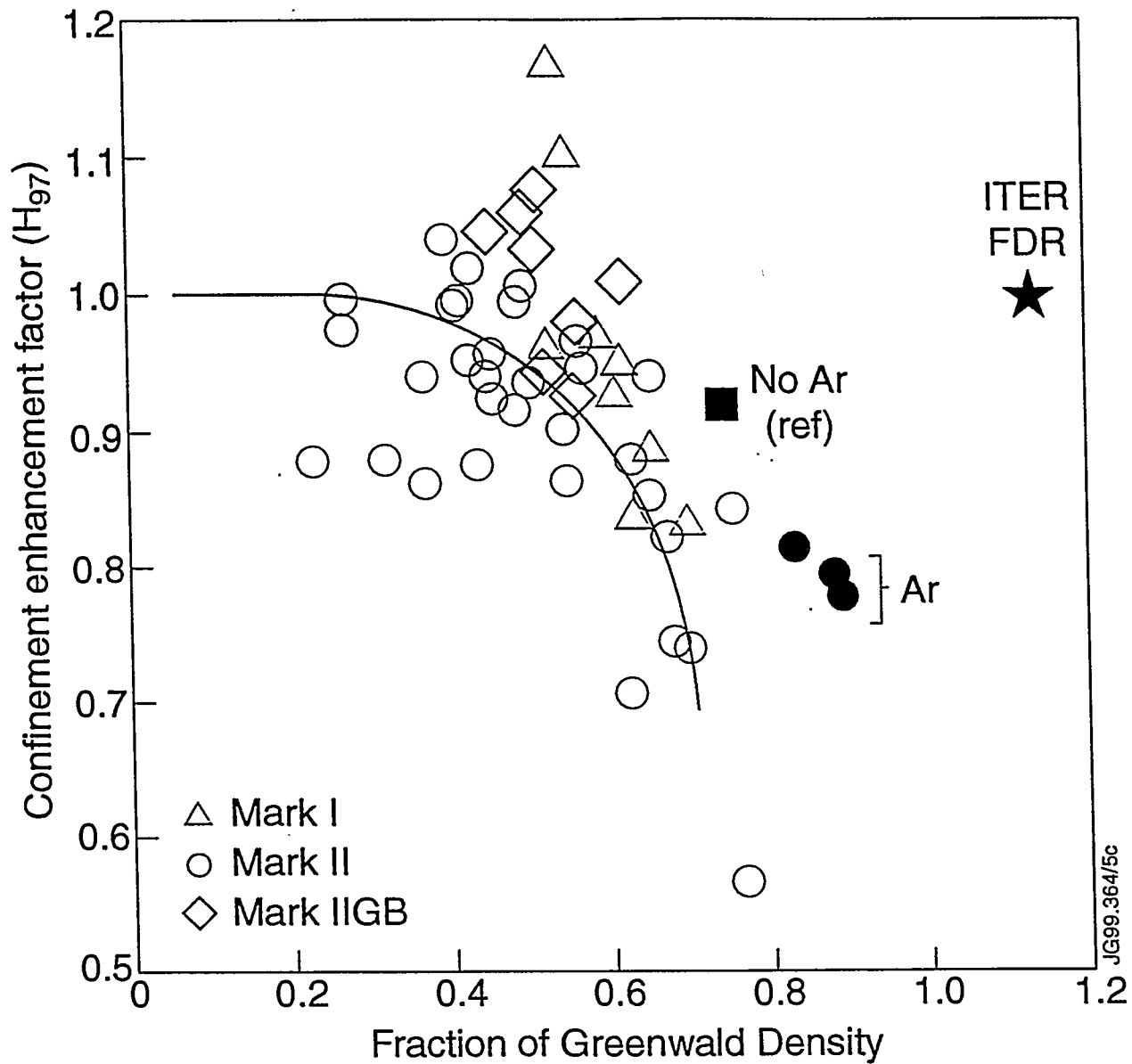


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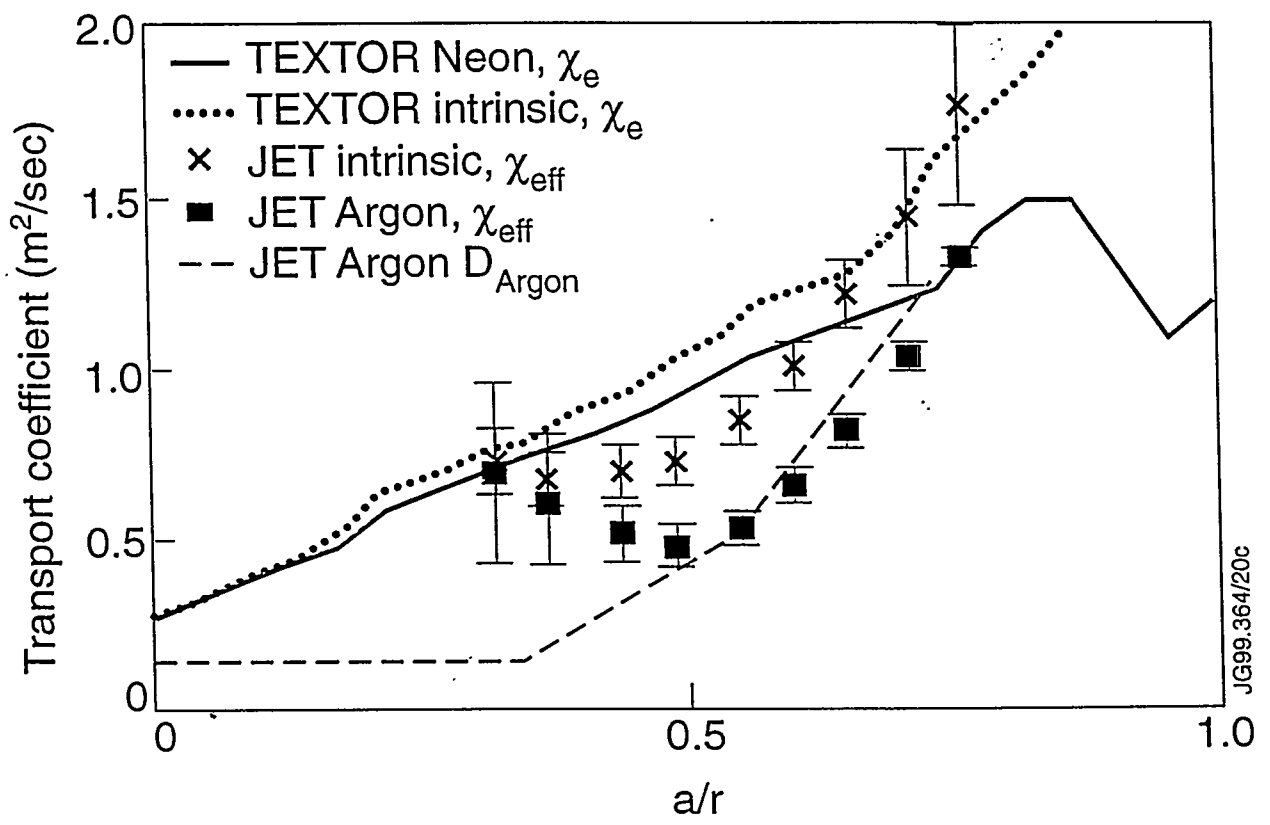


JG99.372/8c

$0.175 < \delta < 0.225$ ,  $2.3 < I_p(\text{MA}) < 2.6$ ,  $2.3 < B_T(\text{T}) < 2.6$ ,  $9 < P_{\text{IN}}(\text{MW}) < 14$







JG99.364/20c

