

BN-600 HYBRID CORE BENCHMARK ANALYSES

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ABSTRACT

Benchmark analyses for the hybrid BN-600 reactor that contains three uranium enrichment zones and one plutonium zone in the core, have been performed within the frame of an IAEA sponsored Co-ordinated Research Project. The results for several relevant reactivity parameters obtained by the participants with their own state-of-the-art basic data and codes, were compared in terms of calculational uncertainty, and their effects on the ULOF transient behavior of the hybrid BN-600 core were evaluated.

The comparison of the diffusion and transport results obtained for the homogeneous representation generally shows good agreement for most parameters between the RZ and HEX-Z models. The burnup effect and the heterogeneity effect on most reactivity parameters also show good agreement for the HEX-Z diffusion and transport theory results. A large difference noticed for the sodium and steel density coefficients is mainly due to differences in the spatial coefficient predictions for non-fuelled regions. The burnup reactivity loss was evaluated to be 0.025 (4.3 \$) within ~ 5.0% standard deviation. The heterogeneity effect on most reactivity coefficients was estimated to be small. The heterogeneity treatment reduced the control rod worth by 2.3%. The heterogeneity effect on the k_{eff} and control rod worth appeared to differ strongly depending on the heterogeneity treatment method. A substantial spread noticed for several reactivity coefficients did not give a significant impact on the transient behavior prediction. This result is attributable to compensating effects between several reactivity effects and the specific design of the partially MOX fuelled hybrid core.

1. INTRODUCTION

This paper presents the results of benchmark analyses of a hybrid UOX/MOX fuelled core of the BN-600 reactor. These benchmark core analyses have been performed within the frame of the IAEA sponsored Co-ordinated Research Project (CRP) on "Updated Codes and Methods to Reduce the Calculational Uncertainties of the LMFR Reactivity Effects," commenced in 1999. The general

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objective of this CRP is to validate, verify and improve methodologies and computer codes used for the calculation of reactivity coefficients in fast reactors aiming at enhancing the utilization of plutonium and minor actinides.

There has been no change in the view that energy production with breeding of fissile materials is the main goal of fast reactor development to ensure long-term fuel supply. However, before the breeding role of fast reactors is recognized economically, due to the increasingly available low-cost uranium from the 1980s onwards, the emphasis of fast reactor development shifted to incineration of stock-piled plutonium and partitioning and transmutation (P&T) of nuclear wastes to meet contemporary demands.

Following a proposal by the Russian Federation, a hybrid UOX/MOX (mixed oxide) fuelled BN-600 reactor core that has a combination of highly enriched uranium (HEU) and mixed oxide (MOX) assemblies in the core region, was chosen as a calculational model [1]. Hence the benchmark clearly addresses the issues of weapons-grade plutonium utilization for energy production in a mixed UOX/MOX fuelled core of the BN-600 reactor.

Nine organizations from eight Member States and the IAEA participated in the hybrid BN-600 core benchmark analyses. The input data for the benchmark neutronics calculations prepared by OKBM and IPPE, have been reviewed and modified in the first meeting of the CRP [1]. The benchmark analyses consist of three Phases during 1999 – 2001: RZ homogeneous benchmark (Phase 1), HEX-Z homogeneous benchmark (Phase 2), and HEX-Z heterogeneous and burnup benchmark (Phase 3).

This paper first addresses the benchmark definitions and specifications given for each Phase and briefly introduces the basic data, computer codes, and methodologies applied to the benchmark analyses by the various participants. Then, the results obtained by the participants in terms of calculational uncertainty and their effect on the core transient behavior are inter-compared. Finally it addresses some preliminary conclusions to be drawn at this stage of the benchmark.

2. DESCRIPTIONS OF BENCHMARKS

2.1. BENCHMARK MODELS

The input data including RZ and HEX-Z calculational models for the BN-600 benchmark calculations are completely described in the benchmark definitions [1, 2]. The hexagonal-Z (HEX-Z) calculational model, a 60° sector of the layout of the benchmark core, is shown in Fig. 1. The calculational model corresponds to the 1470 MW_{th} total power BN-600 reactor at the beginning of an equilibrium cycle, when the impact of the control rods is the strongest. The core consists of a low enrichment inner zone (LEZ), a middle enrichment zone (MEZ), and a high enrichment outer zone (HEZ). Between MEZ and HEZ a mixed oxide zone (MOX) is located. Three control rod (SHR) zones and one scram rod (SCR) zone, consisting of 19 shim control rods and 6 scram control rods respectively, are interspersed radially in LEZ. The outer core zone is bounded by two steel shielding zones, followed by a radial reflector zone. The RZ calculational model indicating the axial positions of the control rods is shown in Fig. 2. In the SHR zone the bottom of the absorber is parked at the core midplane, whereas in the SCR zone the absorber is parked above the enriched fuel region.

From the composition specifications, each enriched fuel zone with axially uniform atomic fractions has an average fuel burnup of 2 – 3%. The LEZ, MEZ, and HEZ have uranium enrichments of 17, 21 and 26wt.% for feed fuel, respectively. The MOX contains plutonium of enrichment 21.3wt.% with

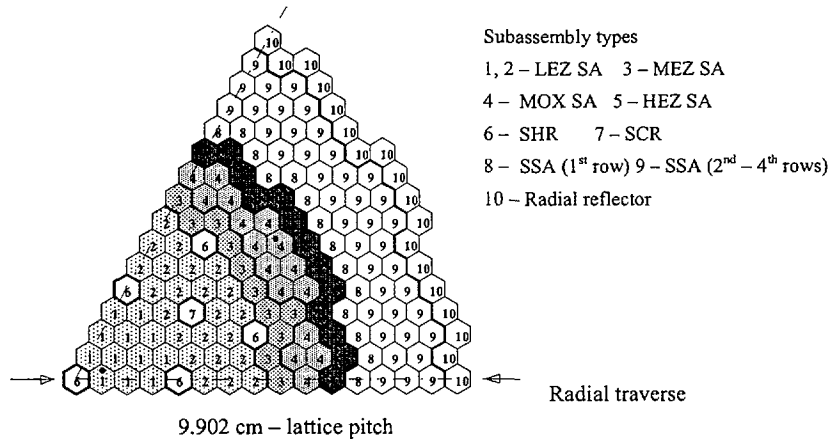


Figure 1. Benchmark model layout (60° sector, rotational symmetry)

DZ	Axial reflector												
20.0	Axial reflector												
24.7	UAB	UAB	SCR	UAB	UAB	UAB	UAB	UAB	UAB				
5.5	SHR	SHR				SHR							
52.2	LEZ-1	LEZ-2	Follow er + rod tail	LEZ-3	SHR	MEZ	MOX	HEZ	SSA	SSA	Radial reflector		
52.2	Follow er + rod tail	LEZ-1	Follow er + rod tail	LEZ-2	Follow er	LEZ-3	Follow er + rod tail	MEZ	MOX	HEZ			
5.5													
29.7	LAB	LAB	LAB	LAB	LAB	LAB	LAB	LAB					
20.0	Axial reflector												
DR (cm)	5.2	36.8	1.8	5.8	1.6	18.9	2.25	10.38	13.5	9.48	9.48	25.43	50.0

Figure 2. RZ calculational model indicating the axial position of control rods

93.8wt.% ²³⁹Pu content for feed fuel. In these specifications, it has been assumed that one fission produces one fission product (FP), corresponding to the definition of FP cross-section in nuclear calculations. All fuel isotopes have been modeled at a uniform temperature of 1500 K and all structural and coolant isotopes are at a uniform temperature of 600 K.

In the Phase 1 and 2 studies the effective multiplication factor (k_{eff}), and the integral value and/or its spatial distribution of reactivity coefficients including kinetics parameters were calculated for the RZ and HEX-Z calculational models by diffusion and transport theory methods, using homogeneous representations of the material regions. Spatial distributions of reactivity coefficients were obtained by first order perturbation theory method. The core power distribution was normalized to the total power of 1470 MW_{th} assuming energy is deposited at the point of fission with an energy of 200 MeV per fission and 0 MeV per capture for all nuclides.

Phase 3 calculations were performed by diffusion and transport theory methods for the HEX-Z model only. For the burnup analysis a single stage calculation has been assumed with no recalculation of the

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flux or resonance self-shielding for sub time steps. The burnup period is 140 effective full power days at an assumed 100% load factor. To evaluate the heterogeneity effect and the burnup effect, k_{eff} , fuel Doppler coefficient and sodium density coefficient were mainly calculated at the beginning of cycle (BOC) and the end of cycle (EOC) for both homogeneous and heterogeneous core models. Detailed heterogeneous geometry configurations were used for the core fuel regions and the control rods in the heterogeneous core model [3]. The control rod worth at BOC and the reactivity loss with burnup were also evaluated.

2.2.2 BASIC DATA AND COMPUTER CODES

Table I summarizes the basic data and computer codes used for the benchmark analyses by the participants, details of which are described in References [1-3]. The multi-group cross-sections were generated based on their own evaluated nuclear data library, using their own current state-of-the-art data processing system. Basically common delayed neutron data sets were used for the calculation of system kinetics parameters. The effective cross-sections condensed into broader energy group structure were generated using their own cell calculation code with various self-shielding treatments for homogeneous and heterogeneous cell models. The energy group number of condensed group constant sets ranges from 230 (fine) to 9 (coarse) with the upper energy boundary of ~ 10 MeV. Various specific heterogeneous cell modeling methods and specific code process schemes were employed for the treatment of the heterogeneity of control rods.

For HEX-Z diffusion theory calculations that are basic calculations in the benchmark analyses, mostly a finite difference method or a nodal diffusion approximation method were employed with the axial mesh size ranging from 2.6 cm to 5.0 cm for core and blanket regions. Most reactivity coefficients were determined using the diffusion theory based, first order perturbation theory method.

3. RESULTS AND DISCUSSIONS

3.1 HOMOGENEOUS BENCHMARKS

3.1.1 Reactivity coefficients and kinetics parameters

The reactivity coefficients for sodium density (W_{Na}), fuel density (W_{fuel}), steel (structure) density (W_{steel}), absorber worth (W_{abs}), and Doppler effect were calculated using the first-order perturbation theory. The density coefficients (material worths) for sodium, fuel and steel density coefficients are defined in terms of $\partial k/k \cdot \rho/\rho$; this is the reactivity change per 1% reduction in material density specified in the benchmark definitions. The fuel and steel Doppler coefficients (K_D^{fuel} and K_D^{steel}) were calculated for temperature changes in the fuel (1500 K to 2100 K) and structure (600 K to 900 K) respectively for a flooded state. The radial and axial expansion coefficients (R_{rad} and R_{ax}) were calculated for a 1% uniform expansion of fuel, steel and absorber within the model in axial and radial directions. They were obtained by the k_{eff} differences using the expanded and un-expanded geometric models. In these definitions, isotope masses are conserved except for sodium, which is kept at constant number density.

Integral reactivity coefficients obtained by the participants for the Phase 1 and 2 benchmarks are given in Tables II and III respectively, with the mean value and relative standard deviation of the results for each coefficient. The comparison of the diffusion and transport results for Phase 1 and 2 generally shows good agreement for most parameters between the RZ and HEX-Z models. The RZ

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and HEX-Z diffusion theory calculations predicted the k_{eff} as 1.00270 and 1.00451 within 0.71% standard deviation, respectively. For the Doppler effect, the structure Doppler was evaluated to be $\sim 1/6$ the magnitude of the fuel Doppler with standard deviations of 12 \sim 25%. These results are similar for Phases 1 and 2.

Table I. Basic nuclear data and computer codes used for benchmark analysis

Participant	ANL	CEA/SA	CIAE	IGCAR	IPPE	JNC	KAERI	OKBM
Nuclear data								
• Nuclear data library	(ENDF/B-V.2)	(JEF-2.2)	LIB-IV-M	CV2M/ ABBN-93	ABBN-93/93.1	JENDL-3.2	KAFAX/ F22	ABBN-93/ 78
• Standard group constant set (No. of groups)	2082	1968	46	25/26	299/26	70	80	26
Delayed neutron data								
• Delayed neutron yield	ENDF/B-V	R.J. Tuttle[4]	R.J. Tuttle	R.J. Tuttle	R.J. Tuttle	R.J. Tuttle	ENDF/B-VI	ABBN-78 (M.C. Brady[5])
• Yield fraction and decay constant	ENDF/B-V	JEF-2.2 (Pu ²⁴² from ENDF/B-VI)	R.J. Tuttle	R.J. Tuttle	R.J. Tuttle	G.R. Keepin[6]	ENDF/B-VI	ABBN-78 (Pu ²⁴¹ , Pu ²⁴² from Tuttle)
• Delayed neutron spectrum	ENDF/B-V	JEF-2.2 (U ²³⁸ , Pu ²⁴² from ENDF/B-VI)	D. Saphier[7]		D. Saphier	D. Saphier	ENDF/B-VI	D. Saphier
Cell calculation								
• Effective cross-section calculation	MC ² -2 ultra-fine B _g ² method	ECCO subgroup method	IDX ABBN f-factor method	EFFCROSS collision prob. method	CONSYST B ² method	SLAROM ABBN f-factor method	TRANSX ABBN f-factor method	CONSYST (for RZ) MIM (for HEX-Z)
• Fuel SA heterogeneity		ECCO	IDX Bell's approx.	NCELL interface current method		CASUP Tone's method	TRANSX Bell's approx.	
• SHR heterogeneity		ECCO/ BISTRO reactivity equivalence method		COHINT collision prob. method		CASUP reaction-rate preservation method	TRANSX	
Core calculation								
<i>Phase 1 (R-Z model)</i>								
• Condensed group const. set (No. of groups)	230	33	12	25/26	18	18	9	26
• Diffusion theory cal.	DIF3D nodal	BISTRO finite difference	2DB finite difference	ALCIALMI finite difference	RHEIN finite difference	CITATION finite difference		SYNTES
• Transport theory cal.	TWODANT S _n method	BISTRO S _n method	DOT3.5 S _n method	DOT S _n method	TWODANT S _n method	TWOTRAN S _n method	TWODANT S _n method	
• Perturbation theory cal.	VARI3D	BISTRO	PERT-V	NEWPERT	RHEIN	PERKY		
• Monte Carlo theory cal.	MCNP		MCNP					
<i>Phases 2&3 (HEX-Z model)</i>								
• Condensed group const. set (No. of groups)	230	33	12	25	18	18	9	26
• Diffusion theory flux/burnup cal.	DIF3D/ REBUS-3 nodal	H3D module TRI-Z, finite difference	HND nodal	3DB/FARC OBAB finite difference	TRIGEX nodal	CITATION finite difference	DIF3D/ REBUS-3 nodal	JARFR finite difference
• Transport theory cal.	VARI3D	TGV/ VARIANT		DOT	TRIGEX	NSHEX nodal, S _n method	SOLTRAN SP ₂ nodal	
• Perturbation theory cal.	VARI3D	H3D	HND	3DPERT		PERKY	PERT-K	
• Monte Carlo theory cal.	MCNP		MCNP					

Table II. Reactivity parameters (RZ model, Phase 1)

Participant	k_{eff}		K_D^{fuel}		K_D^{steel}		W_{Na}		W_{steel}		W_{fuel}		W_{abs}		R_{ex}		R_{rad}	
	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport
ANL	0.99684	1.00788	-0.00652		-0.00110		0.01749		-0.0384		0.3343		-0.0266		0.1327		0.4360	
CEA/SA	1.01676	1.02298	-0.00679	-0.00671	-0.00134	-0.00146	0.00519	0.00343	-0.0112	-0.0147	0.3423	0.3382	-0.0273	-0.0263	0.1378	0.1374	0.4650	0.4573
CIAE	0.99808	1.01497	-0.00499		-0.00049		0.00211		-0.0055		0.3492		-0.0391		0.1432	0.1341	0.4232	0.4603
IGCAR	1.00360		-0.00462				0.00446		-0.0021		0.3410		-0.0269		0.1395		0.4840	
IPPE	1.00140	1.00576	-0.00622	-0.00628	-0.00122	-0.00124	0.00898	0.00201	-0.0053	-0.0088	0.3505	0.3466	-0.0274	-0.0260	0.1297	0.1267	0.4868	0.4787
JNC	1.00423	1.00948	-0.00635	-0.00621	-0.00116	-0.00113	0.00765	0.00403	-0.0126	-0.0149	0.3491	0.3491	-0.0265	-0.0270	0.1393	0.1361	0.4812	0.4647
KAERI		1.02654		-0.00777		-0.00125										0.1376		0.4493
OKBM	0.99796		-0.00659		-0.00097		0.01065		-0.0114		0.3628		-0.0287		0.1415		0.4827	
Mean	1.00270	1.01460	-0.00601	-0.00674	-0.00105	-0.00127	0.00808	0.00319	-0.0123	-0.0128	0.3470	0.3446	-0.0287	-0.0264	0.1377	0.1344	0.4656	0.4621
SD (\pm)	0.00623	0.00777	0.00079	0.00062	0.00027	0.00012	0.00468	0.00080	0.0112	0.0028	0.0084	0.0047	0.0043	0.0004	0.0045	0.0040	0.0239	0.0097
(Rel. %)	(0.62)	(0.77)	(13.1)	(9.2)	(25.7)	(9.4)	(57.9)	(25.1)	(91.1)	(21.9)	(2.4)	(1.4)	(15.0)	(1.5)	(3.3)	(3.0)	(5.1)	(2.1)

Table III. Reactivity parameters (HEX-Z model, Phase 2)

Participant	k_{eff}		K_D^{fuel}		K_D^{steel}		W_{Na}		W_{steel}		W_{fuel}		W_{abs}		R_{ex}		R_{rad}	
	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport	Diffusion	Transport
ANL	0.99592	0.99802	-0.00665		-0.00106		0.02072		-0.0284		0.3339		-0.0241		0.1466		0.4427	
CEA/SA	1.01546	1.02272	-0.00686	-0.00681	-0.00131	-0.00132	0.00771	0.00592	-0.0022	-0.0073	0.3435	0.3385	-0.0258	-0.0234	0.1463	0.1397	0.4741	0.4605
CIAE	1.00026		-0.00610		-0.00067		0.00486		-0.0037		0.3591		-0.0243		0.1600		0.5008	
IGCAR	1.00642		-0.00471				0.00668		0.0107		0.3691		-0.0234		0.1461		0.4788	
IPPE	1.00471		-0.00635		-0.00080		0.00877		0.0019		0.3477		-0.0217		0.1384		0.4904	
JNC	1.00713	1.00967	-0.00646	-0.00633	-0.00109	-0.00109	0.01038	0.00664	-0.0020	-0.0066	0.3466	0.3457	-0.0200	-0.0219	0.1574	0.1530	0.4923	0.4814
KAERI	1.01260	1.01494	-0.00770	-0.00764	-0.00108	-0.00106	0.01050		0.0394		0.3490		-0.0267		0.1528	0.1514	0.4726	0.4679
OKBM	0.99357		-0.00632		-0.00106		0.01226		-0.0019		0.3561		-0.0245		0.1293		0.4941	
Mean	1.00451	1.01134	-0.00639	-0.00693	-0.00101	-0.00116	0.01027	0.00628	-0.0017	-0.0069	0.3506	0.3421	-0.0238	-0.0227	0.1471	0.1480	0.4807	0.4699
SD (\pm)	0.00714	0.00898	0.00078	0.00054	0.00019	0.00012	0.00453	0.00036	0.0176	0.0004	0.0102	0.0036	0.0020	0.0008	0.0094	0.0059	0.0172	0.0087
(Rel. %)	(0.71)	(0.89)	(12.2)	(7.8)	(18.8)	(10.3)	(44.1)	(5.7)	(1354)	(58.0)	(2.9)	(1.1)	(8.4)	(3.5)	(6.4)	(4.0)	(3.6)	(1.9)

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The fuel density coefficient was estimated to be ~ 30 times larger than the steel density coefficient in the Phase 1 diffusion results. The fuel density coefficient was estimated even greater times larger for Phase 2. The reactivity coefficient distributions show that the fuel density coefficient is positive and the steel density is negative in the core zone, and that their signs are reversed in the blanket zone. The sodium density coefficients are 0.0081 and 0.0032 for the Phase 1 diffusion and transport theory results. The sodium density coefficients were evaluated to be larger for Phase 2. The absorber worth is ~ -0.028 and ~ -0.023 for Phases 1 and 2. The radial and axial expansion coefficients were evaluated to be ~ 0.14 and ~ 0.47 for Phases 1 and 2. The radial expansion coefficient is ~ 3 times larger than the axial expansion coefficient for a given relative change in dimension.

In the Phase 1 results, the sodium and steel density coefficients were obtained with standard deviations of 58% and 91% respectively in the diffusion theory results and also have a relatively large difference between the diffusion and transport theory results. These results are very similar for Phase 2. In particular, the steel density coefficient has a considerable dispersion. These results indicate that the transport theory method is more appropriate for the calculation of the sodium and steel density coefficients. In addition, the results for the steel density coefficient, even obtained by the transport theory method, show a considerable difference between the Phase 1 and 2 results. This is mainly attributed to differences in the spatial coefficient predictions for non-fuelled regions.

The effective delayed neutron fraction (β_{eff}) and prompt neutron lifetime (l_p) for system kinetics parameters are given in Table IV. The results for Phases 1 and 2 show good agreement between the RZ and HEX-Z models and are very similar. The obtained effective delayed neutron fraction, $\beta_{eff} = 0.0059$ is representative of a fast reactor system dominated by ^{235}U fission ($\beta = 0.0064$), but with the contribution from ^{239}Pu ($\beta = 0.0020$).

Table IV. Kinetics parameters

Participant	RZ Model (Phase 1)		HEX-Z Model (Phase 2)	
	β_{eff} (pcm)	l_p (10^{-7} sec)	β_{eff} (pcm)	l_p (10^{-7} sec)
ANL	569	4.423	581	4.491
CEA/SA	597 (594)*	4.687	598	4.606
CIAE	604	4.277	598	4.239
IGCAR	598	4.618	598	
IPPE	585 (583)*	4.586	593	
JNC	586	4.620	587	4.484
KAERI			595	4.443
OKBM	578	4.618	573	4.130
Mean	588	4.514	590	4.399
SD (\pm) (Rel. %)	11.2 (1.9)	0.134 (3.0)	8.9 (1.5)	0.163 (3.7)

* Transport theory results

The regionwise sodium density coefficients for the core plus the upper axial blanket (UAB) regions obtained from the RZ and HEX-Z diffusion calculations are shown in Fig. 3, with the sum of the entire regions for comparison. The integral comparisons using only the core region plus the upper

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axial blanket regions are given for a better gauge of comparison, since these regions will experience significantly higher temperatures in an accident situation. Furthermore, these regions are the locations where the largest cross-section differences would be expected. From the comparison of these figures, it can be inferred that large differences between the values obtained by the participants are due to the coefficients from the control regions and other non-fuelled regions. In these figures, the comparison of regional coefficients shows better agreement than the comparison of total coefficients summed over all regions, because the total value is a sum of very small positive and small negative contributions calculated from a 1% reduction in the material density. The resultant aspects are very similar for the steel density coefficients shown in Fig. 4.

The sodium density coefficient was investigated by separating its regionwise distributions for the leakage and non-leakage components. The regionwise leakage and non-leakage components for the core fuelled regions obtained from HEX-Z diffusion calculations are also shown in Fig. 3 with the averaged values. The comparison of the regionwise components shows similar aspects noticed for the sodium and steel density coefficients, particularly in the MEZ and MOX regions. In these figures, it is clearly observed that the leakage component becomes dominant in HEZ, and, as a result, HEZ has the greatest positive coefficient.

3.1.2 Power distributions

The core powers were calculated based on a local energy deposition model where energy is deposited at the point of fission with an energy of 200 MeV per fission and 0 MeV per capture for all nuclides. This local energy deposition model does not allow for gamma transport toward the outer regions of the core, and, as a result, the powers for non-fuelled regions are completely ignored.

Regionwise normalized power distributions calculated by the RZ and HEX-Z diffusion calculations are shown in Fig. 5 with their average values. The comparison of regionwise power distributions show good agreement for the RZ and HEX-Z models.

3.1.3. Influence of mesh size and energy group numbers

As previously mentioned, various planar and axial mesh sizes were used in the benchmark calculations, even though the axial mesh sizes were specified for desired result edit. In addition, some results were obtained using relatively low group numbered cross-sections in few group calculations. The value of a corresponding additional uncertainty strongly depends on the number of energy groups and the choice of the energy group boundaries. The OKBM participant devoted to investigate the possible uncertainties due to coarse mesh and few group calculations of the reactivity coefficients [8]. This parametric study was performed based on the Phase 2 results obtained using the finite difference option of the JARFR code in HEX-Z/triangular-Z diffusion approximations. In this study, few group calculations were performed for 26 to 9 energy groups with axial mesh sizes being varied from 2.61 cm to 10.0 cm for the core and from 2.5 cm to 5.0 cm for the axial blankets, respectively.

Among the examined parameters, the planar mesh size appeared to be most influential. The HEX-Z option has a coarse planar mesh configuration, i.e., one mesh point per hexagon, compared with six points per hexagon of the triangular-Z option. The results show that the use of the HEX-Z option having a larger planar mesh, results in the following additional uncertainties; overestimation of the fuel and steel Doppler coefficients by up to ~10 – 15% in the upper blanket, some shift of the sodium worth to positive by up to ~ 20% in LEZ, and significant underestimation of the absorber worth by 20%. The use of a larger axial mesh leads to some overestimation by 4 – 8% of all reactivity coefficients in axial blankets.

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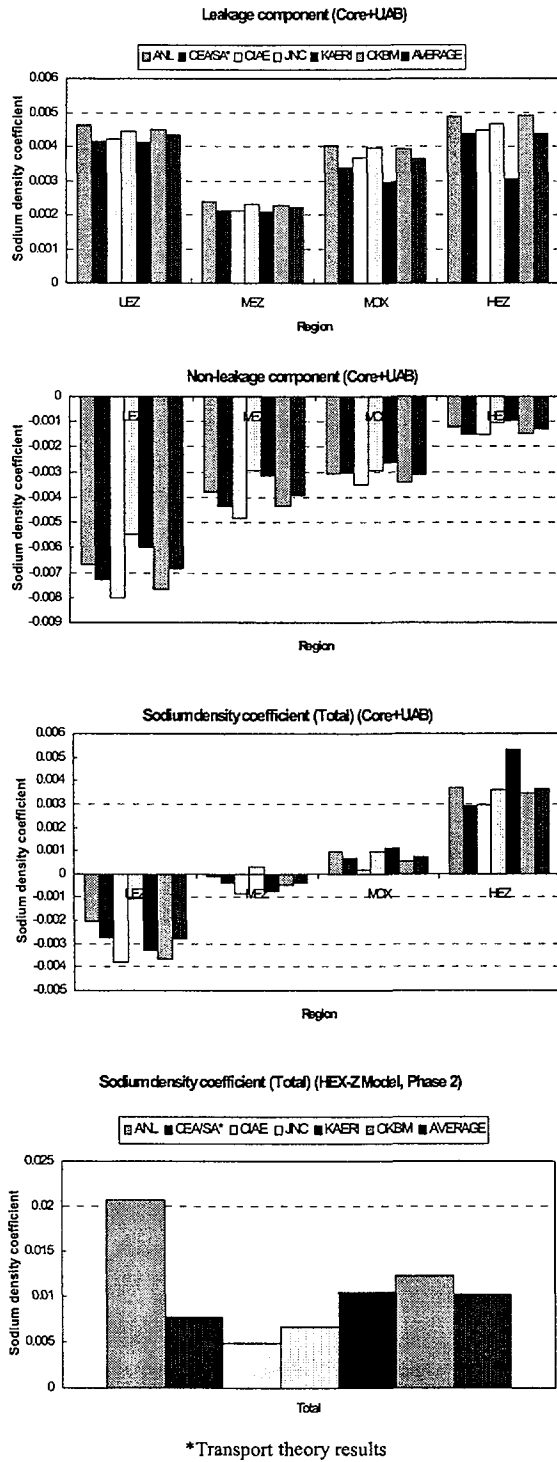


Figure 3. Sodium density coefficients

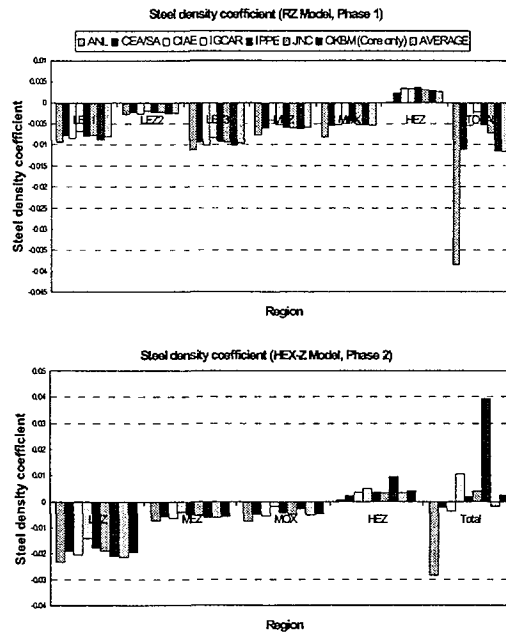
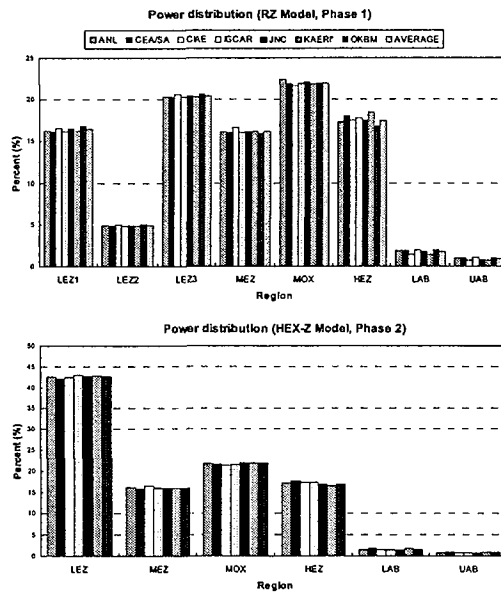


Figure 4. Steel density coefficients



LEZ = LEZ1 + LEZ2 + LEZ3 in the RZ model

Figure 5. Power distributions

It has also been revealed that the few group schemes, used by some participants, may result in significant uncertainties by up to -45% in the core and by up to -24% in the axial blankets for the sodium worth. The principle of the choice of the energy scale optimal division has been proposed for the few group calculations. Based on this principle, the optimal collapsing of the initial 26 groups into 9 groups has been provided. The sodium worth in the core (except for the MOX core), obtained from the 9 group calculation, agreed well with that obtained from 26 group calculation.

3.2 HETEROGENEOUS AND BURNUP BENCHMARKS

3.2.1 Burnup effect

To evaluate the burnup effect, k_{eff} , fuel Doppler coefficient and sodium density coefficient were calculated at the beginning of cycle (BOC) with the SHR control rods mid-core insertion and the end of cycle (EOC) with the SHR rods fully withdrawn for the homogeneous core model. Burnup calculations were performed for 140 effective full power days at 1470 MW_{th} in the equilibrium cycle with SHR control rods mid-core insertion. The k_{eff} at EOC with the SHR rods mid-core position was calculated to evaluate the reactivity loss with burnup. The control rod worth for the whole bank of rods were calculated for the homogeneous model, one with the SHR rods fully "withdrawn" and the other with the rods fully "inserted". Core power calculations assumed a local energy deposition model using an energy 200 MeV per fission for fissionable isotopes and 0 MeV per capture for all nuclides.

The reactivity parameters at BOC and EOC obtained by the homogeneous core model are given in Table V. Most burnup effects on the reactivity parameters show good agreement for diffusion and transport theory results. The reactivity loss with burnup in an equilibrium cycle was predicted to be 0.025 (4.3 \$) both from the diffusion and transport theory results within 5.0% standard deviation. The fuel Doppler coefficient is more negative by ~ 6.0% and the absorber worth is less negative by ~ 3.0% at EOC, respectively. This is counterintuitive since the ²³⁸U content in the core is reduced at EOC. It is conjectured that this change is caused by the power shift toward the LEZ zone. The burnup effects on the steel Doppler coefficient appeared to be negligibly small.

The sodium density coefficient obtained for the whole core becomes more positive at EOC by 53% at maximum. This is attributed to the build up of fission products leading to a harder spectrum, and is exacerbated by the power shift into the low leakage LEZ zone. However, in details, the sodium worth is observed to be less positive in most results. This conflicting result is surmised to come from the summation over its very small spatial contributions from non-fuelled regions, particularly at EOC, similar for Phases 1 and 2. The conflicting results cause a large standard deviation in the determination of the sodium worth. The steel density coefficient seems to have no difference in its value between BOC and EOC. The absorber worth becomes less negative by 3.5% at EOC. The control rod worth is more positive by 2.5% at EOC, and its burnup effect is small.

3.2.2 Heterogeneity effects

The k_{eff} , control rod worth and sodium density coefficient were calculated with the heterogeneous treatment of the core fuel regions (LEZ, MEZ, HEZ, and MOX, and the core and axial blankets) and the SHR control rods in the heterogeneous core model. Several insignificant simplifications of real geometry description have been assumed in the heterogeneous geometry descriptions Fuel subassemblies for LEZ, MEZ, MOX, HEZ have the identical geometry, where 127 fuel pins are located with triangular pitch of 7.95 mm inside hexagonal wrapper. A separate geometry description was defined for the SHR absorber region [3].

Table V. Burnup effects

Participant	k_{eff}			K_D^{fuel}			K_D^{steel}			W_{Na}		
	BOC	EOC	Bumup reactivity loss ¹⁾	BOC	EOC	Bumup effect (%) ²⁾	BOC	EOC	Bumup effect (%)	BOC	EOC	Bumup effect (%)
ANL	0.99592	0.97162	0.02511	-0.00665	-0.00687	3.37	-0.00106			[0.00080] ⁴⁾		
CEA/SA	1.01546 (1.02272) ³⁾	0.98914 (0.99642)	0.02621 (0.02581)	-0.00683 (-0.00682)	-0.00732 (-0.00724)	7.11 (6.15)	-0.00128 (-0.00128)	-0.00129 (-0.00129)	1.25 (-1.90)	0.00610 (0.00592)	0.00943 (0.00762)	54.53 (24.78)
IGCAR	1.00136 (1.00188)	0.97671	0.02520	-0.00684	-0.00731	6.87				0.01090	0.00920	-15.60
JNC	1.00713 (1.00967)	0.98323 (0.98542)	0.02414 (0.02438)	-0.00646 (-0.00633)	-0.00703 (-0.00692)	8.82 (9.32)	-0.00109 (-0.00109)			0.01038 (0.00664)	0.00937 (0.00542)	-9.73 (-18.37)
KAERI	1.01467	0.99051	0.02404	-0.00770 (-0.00764)	-0.00777	0.98	-0.00108 (-0.00106)			0.01050	0.00831	-20.80
OKBM	0.99486	0.96815	0.02773	-0.00632	-0.00682	9.82	-0.00106			0.01226	0.01105	-9.87
Mean	1.00430 (1.01014)	0.97989 (0.99091)	0.02540 (0.02510)	-0.00666 (-0.0693)	-0.00719 (-0.00708)	6.16 (7.73)	-0.00210 (-0.00116)	-0.00129 (-0.00129)	1.25 (-1.90)	0.00911 (0.00628)	0.00913 (0.00652)	4.83 (3.21)
SD (±)	0.00729 (0.01009)	0.00843 (0.00550)	0.00139 (0.00101)	0.00046 (0.00054)	0.00032 (0.00016)	3.36 (2.24)	0.00262 (0.00011)	0.00	0.00 (0.00)	0.00254 (0.00036)	0.00449 (0.00110)	28.79 (30.51)
Participant	W_{steel}			W_{fuel}			W_{abs}			Control rod worth		
	BOC	EOC	Bumup effect (%)	BOC	EOC	Bumup effect (%)	BOC	EOC	Bumup effect (%)	BOC	EOC	Bumup effect (%)
CEA/SA	-0.0041 (-0.0088)	-0.0041 (-0.0088)	0.00 (0.00)	0.3344 (0.3385)	0.3431 (0.3369)	-0.40 (-0.46)	-0.0275 (-0.0234)	-0.0265 (-0.0228)	-3.54 (-2.60)	0.0697 (0.0652)	0.0697 (0.0652)	0.00 (0.00)
IGCAR	0.0007	-0.0087	-1342.86	0.3691						0.0682	0.0716	4.91
Mean	-0.0078 (-0.0077)	-0.0064 (-0.0088)	-671.43 (0.00)	0.3507 (0.3421)	0.3495 (0.3369)	-0.63 (-0.46)	-0.0235 (-0.0226)	-0.0265 (-0.0228)	-3.54 (-2.60)	0.0661 (0.0652)	0.0706 (0.0652)	2.46 (0.00)
SD (±)	0.0123 (0.0011)	0.0023 (0.00)	0.00 (0.00)	0.0096 (0.0036)	0.0065 (0.00)	0.32 (0.00)	0.0025 (0.0007)	0.00 (0.00)	0.00 (0.00)	0.0029 (0.00)	0.0009 (0.00)	3.47 (0.00)

Notes: 1) Burnup reactivity loss = $(k_{eff}^{BOC} - k_{eff}^{EOC}) / k_{eff}^{BOC} / k_{eff}^{EOC}$ based on diffusion results, and transport results denoted by parenthesis.

2) Burnup effect = $(BOC - EOC) / BOC \times 100$ (%) based on diffusion results, and transport results denoted by parenthesis.

3) Values in parentheses denote transport results.

4) Value only for the core and upper axial blanket regions.

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To evaluate the heterogeneity effect, the heterogeneous geometries of the SHR control rods have been treated by employing various specific heterogeneous modeling methods with specific procedures for reaction rate preservation [9-12]. The control rod worth for the whole bank of rods were calculated for the heterogeneous model, one with the SHR rods fully "withdrawn" and the other with the rods fully "inserted".

The reactivity parameters calculated by the homogeneous and heterogeneous core models with the SHR rods mid-core insertion at BOC are given in Table VI. Most reactivity parameters and resultant heterogeneity effects show good agreement for diffusion and transport theory results. The heterogeneity effect on k_{eff} at BOC was evaluated to be the increase of 40 – 92 pcm and was observed to be negligibly small. The obtained results appeared to differ depending on the heterogeneity treatment method and, as a result, it leads to a large standard deviation in the heterogeneity prediction. The fuel Doppler coefficient becomes less negative by ~ 6.0% with the heterogeneous treatment. The heterogeneity effect on the steel Doppler coefficient appeared to be negligibly small (based on the transport results). The heterogeneous treatment increases the sodium density coefficient by 53% at maximum. However, there remains a large standard deviation in the prediction between the obtained values.

The heterogeneity effect on the fuel, steel and absorber worth is negligibly small. The control rod worth becomes less positive by 2.3% (2.3% decrease in the control rod worth). From the comparison of the obtained values, it has been observed that the prediction of heterogeneity effect strongly depends on the heterogeneous modeling and the heterogeneity treatment methodology.

3.2.3 Influence of basic nuclear data

As part of the BN-600 hybrid core benchmark analysis effort, the JNC participant performed the sensitivity analysis of the difference between the JENDL-3.2 and JEF-2.2 libraries [13, 14]. Using the cross-section sensitivity coefficients, the nuclide-wise and reaction-wise contribution to the effect of the library change was analyzed for important reactivity parameters. From the sensitivity analysis, dominant nuclides and reactions were identified for those reactivity parameters. For example, the ^{235}U and iron cross-sections largely contribute to the k_{eff} difference between JEF and JENDL. The ^{235}U capture and iron elastic reactions are the largest components of the sodium density reactivity difference and the ^{235}U capture and ^{238}U capture reactions are main contributors to the burnup reactivity loss difference.

Major findings observed from the results of the sensitivity analysis are given as follows: the criticality and sodium density reactivity are very sensitive to cross-section changes, and show strong non-linearity, which makes difficult for precise design work. The steel density reactivity also has large sensitivity coefficients, but the mechanism seems to be rather simple. The fuel and absorber density reactivities are simple and easy to predict.

3.3 ULOF TRANSIENT ANALYSES

The Russian participants, OKBM and IPPE, performed ULOF transient studies using the results prepared by the participants in Phases 1 and 2 [15, 16]. These transient analyses have been proposed to preliminarily evaluate the effect of possible errors occurred in the reactivity coefficient calculations by comparing reactor dynamics behavior under several accident conditions. IPPE performed a simplified transient analysis up to onset of sodium boiling and OKBM carried out SAS4A transient analysis up to beyond sodium boiling with a more detailed reactor model.

Table VI. Heterogeneity effects

Participant	k_{eff}			K_D^{fuel}			K_D^{steel}			W_{Na}		
	Diffusion	Transport	Hete. effect (pcm) ¹⁾	Diffusion	Transport	Hete. effect (%) ²⁾	Diffusion	Transport	Hete. effect (%)	Diffusion	Transport	Hete. effect (%)
ANL	(0.99592) ³⁾	(0.99802)		(-0.00665)			(-0.00106)			[0.00080] ⁴⁾		[0.12]
CEA/SA	1.01474 (1.01546)	1.02211 (1.02272)	-72 (-61)	-0.00734 (-0.00683)	-0.00722 (-0.00682)	7.48 (5.93)	-0.00128 (-0.00128)	-0.00131 (-0.00132)	0.55 (-1.06)	0.01062 (0.00610)	0.00984 (0.00592)	73.99 (66.29)
CIAE	1.00214 (1.00026)		188	-0.00640 (-0.00610)		4.77	-0.00211 (-0.00067)		214.25	0.00529 (0.00487)		8.65
IGCAR	1.00188 (1.00136)		52	(-0.00684)						0.00838 (0.01090)		-23.12
JNC	1.00913 (1.00713)	1.01109 (1.00967)	200 (142)	(-0.00646)	(-0.00633)		(-0.00109)	(-0.00109)		0.01280 (0.01038)	0.00934 (0.00664)	23.31 (40.66)
KAERI	(1.01260)			(-0.00770)	(-0.00764)		(-0.00108)	(-0.00106)		0.00858 (0.01050)		-18.30
Mean	1.00697	1.01660	92 (40)	-0.00664	-0.00722	6.13 (5.93)	-0.00169	-0.00131	107.40 (-1.06)	0.00913	0.00959	12.91 (53.48)
SD (±)	0.00534	0.00551	111 (102)	0.00047	0.00	1.92 (0.00)	0.00041	0.00	151.11 (0.00)	0.00449	0.00025	39.14 (18.12)
Participant	W_{steel}			W_{fuel}			W_{abs}			Control rod worth		
	Diffusion	Transport	Hete. effect (%)	Diffusion	Transport	Hete. effect (%)	Diffusion	Transport	Hete. effect (%)	Diffusion	Transport	Hete. effect (%)
ANL	(-0.0037)			(0.3339)			(-0.0241)			(0.0653)	0.0598 ³⁾	(-4.78)
CEA/SA	0.0027 (-0.0041)	-0.0015 (-0.0088)	-165.44 (-83.49)	0.3431 (0.3444)	0.3369 (0.3385)	-0.40 (-0.46)	-0.0265 (-0.0275)	-0.0228 (-0.0234)	-3.54 (-2.60)	0.0683 (0.0697)	0.0641 (0.0652)	-1.94 (-1.69)
CIAE	(-0.0037)			0.3560 (0.3591)		-0.85	(-0.0243)			0.0664 (0.0662)		0.33
IGCAR	(0.0007)									0.0656 (0.0682)		-3.77
JNC	(-0.0020)	(-0.0066)		(0.3466)	(0.3457)		(-0.0200)	(-0.0219)		0.0579 (0.0603)	0.0621 (0.0653)	-3.98 (-4.90)
Mean	0.0027	-0.0015	-165.44 (-83.49)	0.3495	0.3369	-0.63 (-0.16)	-0.0265	-0.0228	0.10 (0.06)	0.0646	0.0631	-2.34 (-3.30)
SD (±)	0.00	0.00	0.00 (0.00)	0.0065	0.00	0.32 (0.00)	0.00	0.00	0.00 (0.00)	0.0040	0.0010	2.00 (2.27)

- Notes: 1) Heterogeneity effect = (Hete. - Homo.) x 10⁵ (pcm) based on diffusion results, and transport results denoted by parenthesis
 2) Heterogeneity effect = (Hete. - Homo.) / Homo. x 100 (%) based on diffusion results, and transport results denoted by parenthesis
 3) Values in parentheses denote homogeneous results at BOC.
 4) Values in brackets denote only for the core and upper axial blanket regions.
 5) Monte Carlo results.

- MW -

The simplified transient analysis employs point kinetics approximation without taking into account phase transitions of core materials and sodium boiling. The simplified ULOF ($G = 30\% G_0$) transient analysis shows that flow of ULOF is characterized by the reactivity balance mainly due to the fuel Doppler effect and radial expansion in the hybrid BN-600 core. No principal differences in the ULOF flow were found in cases of using the Phase 2 diffusion results of different participants, especially for reactivity coefficients due to thermal expansion of materials. For the hybrid BN-600 core the influence of reactivity due to sodium density change on the accident process appeared to be negligible.

The SAS4A transient analysis used a simplified channel model for the BN-600 reactor hybrid core. Total eight input data sets were prepared for SAS4A code calculations using the reactivity coefficients and power distributions obtained by the diffusion approximation, taking into account the energy release at the point of neutron capture in the Phase 2 study. The results of the inter-comparison of the ULOF transient analyses show that the existing spread in the reactivity coefficients is insignificant from the viewpoint of consequent spread in the basic integrated parameters of the initial ULOF accident stage. The spread in reactivity coefficients correlates with the spread (up to the sign reverse for the sodium density component) of the values of corresponding full current reactivity components.

However, after the onset of sodium boiling, the spread in the reactivity coefficient distributions of the Doppler effect and the sodium worth in associated with the power distributions causes a significant divergence in the characteristics of the accident progress. Variations of the reactivity components obtained based on the OKBM data are illustrated in Figs. 6 and 7. The sodium density component shows the most significant relative spread among the components of the current reactivity balance. Up to the onset of sodium boiling, this large spread causes the current reactivity ranging from -66% to 59%. At the sodium boiling stage, the impact of a large uncertainty in the sodium worth becomes much more substantial. A significant contribution of the components from core radial deformation and the large uncertainty of its value have been noticed.

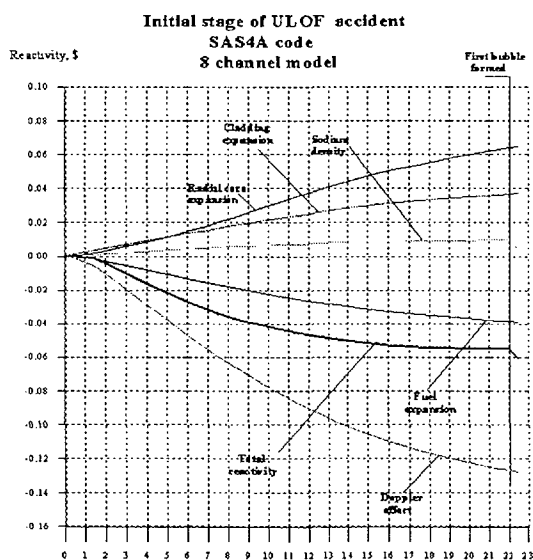


Figure 6. Time dependence of reactivity components before the onset of sodium boiling, based on OKBM data

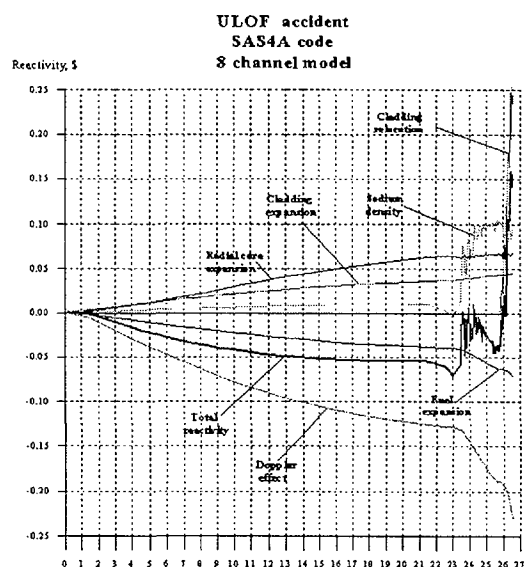


Figure 7. Time dependence of reactivity components beyond sodium boiling, based on OKBM data

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It is noteworthy that the results of the inter-comparison of the ULOF transient analyses using the reactivity coefficients and power distributions of the different participants show good agreement up to the onset of sodium boiling in spite of a substantial spread noticed for several reactivity coefficients. This is due to compensating effects between several reactivity effects and the specific design of the hybrid core mainly loaded with UOX fuel. On the other hand, beyond sodium boiling there exist significant differences in the transient behaviours resulting mainly from differences in Doppler, sodium void and power distributions.

CONCLUSIONS

Benchmark analyses for the UOX/MOX hybrid BN-600 reactor that contains three uranium enrichment zones and one plutonium zone in the core, have been performed within the frame for the IAEA sponsored Co-ordinated Research Project. In these benchmark analyses, results for several important reactivity coefficients obtained by the participants with their own state-of-the-art methodologies, were compared in terms of calculational uncertainties. In addition, the latter's effects on the ULOF transient behaviour of the hybrid BN-600 core were evaluated.

The comparison of the diffusion and transport results for Phase 1 and 2 obtained for the homogeneous representation generally show good agreement for most parameters between the RZ and HEX-Z models. The results for the sodium and steel density coefficients show large discrepancies both in the diffusion approximation, and between the diffusion and transport theory approximations. Large differences in these integral reactivity coefficients result from summation over very small contributions from the control rod regions and other non-fuelled regions. The separate treatment of the leakage and non-leakage components of the sodium density coefficient clearly indicates a similar aspect. In the Phase 3 study, the burnup effect and the heterogeneity effect on most reactivity parameters show good agreement for diffusion and transport theory results. The burnup reactivity loss was predicted to be 0.025 (4.3 \$) within ~ 5% standard deviation. The heterogeneity effect on most reactivity parameters was evaluated to be small. The control rod worth decreases by 2.3% with the heterogeneous treatment. It has been observed that the heterogeneity effect on k_{eff} and control rod worth differs strongly depending on the heterogeneity treatment method.

In the ULOF transient analyses, the substantial spread between the different participants noticed for several reactivity coefficients and power distributions did not have a significant impact on the transient behavior prediction, especially up to the onset of sodium boiling. This result is attributable to compensating effects between several reactivity effects in the specific design of the hybrid core mainly loaded with UOX fuel. This gives confidence that the outcome of this type of transient could be understood in the partially MOX fuelled hybrid core type.

The benchmark analyses will be extended to the study of a BN-600 fully fuelled MOX core design with sodium plenum above the core, and the transient analyses in the fully fuelled MOX core. It remains of significant interest to the participants to seek experimental verification of the reactivity coefficient distributions and integral coefficients through the analysis of experimental measurements.

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